

An r - h Adaptive Kinematic Approach for 3D Limit Analysis

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Abstract

This paper explores a pathway for increasing efficiency in numerical 3D limit analysis through r - h adaptivity, wherein nodal positions (r) and element lengths (h) are successively refined. The approach uses an iterative, nested optimization procedure involving three components: (1) determination of velocities for a fixed mesh of rigid, translational elements (blocks) using second-order cone programming; (2) adaptation of nodal positions using non-linear optimization (r adaptivity); and (3) subdivision of elements based on the magnitude of the velocity jumps (h adaptivity). Examples show that the method can compute reasonably accurate limit loads at relatively low computational cost.

Keywords: limit analysis, 3D, kinematic method, upper bound, adaptivity

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1. Introduction

Accurate evaluation of the limit load, or collapse load, causing failure of a mass of geomaterial is crucial for the design of geotechnical infrastructure, including foundations, slopes, and earth retaining systems. Limit load computations are also central in the determination of how to induce failure deliberately, as in excavation, mining, and earthmoving (e.g., Hettiaratchi and Reece, 1974; Godwin and O’Dogherty, 2007; Hambleton et al., 2014; Hambleton, 2017). Many models rely on a two-dimensional (2D) idealization of the true configuration (e.g., plane strain or axisymmetry), which significantly simplifies the calculations. However, in many cases, the three-dimensional (3D) nature of the problem cannot be ignored. When 3D conditions prevail, computations based on the 2D simplification can overestimate or underestimate the limit load (Soubra and Regenass, 2000; Antão et al., 2011; Michalowski, 2001; Griffiths and Marquez, 2007; Michalowski and Drescher, 2009; Wörden and Achmus, 2013).

Among various existing methods, the kinematic approach of limit analysis is a particularly effective and useful means of evaluating limit loads (cf. Chen, 1975). The kinematic theorem states that for any kinematically admissible velocity field (i.e., failure or collapse mechanism), the load computed by equating the work rate of external forces to the internal energy dissipation rate is a rigorous bound on the true limit load. It gives an upper bound for a load inducing collapse and a lower bound for a load resisting collapse (Drescher, 1991). A kinematically admissible velocity field is one that satisfies boundary conditions and the plastic flow rule. The kinematic theorem requires that material is perfectly plastic and obeys the associative

26 flow rule. The consequences of associativity and possible workarounds in
 27 instances where it may lead to unrealistic predictions are discussed by vari-
 28 ous authors (Davis, 1968; Davis and Booker, 1971; Chen, 1975; Drescher and
 29 Detournay, 1993; Krabbenhoft et al., 2012; Sloan, 2013). For 3D problems
 30 with simple geometries and loading conditions, a kinematically admissible
 31 mechanism can be constructed manually, thereby permitting an analytical
 32 or semi-analytical solution (e.g., Murray and Geddes 1987; Soubra and Re-
 33 genass 2000; Michalowski 2001). Nevertheless, it is generally difficult to
 34 construct collapse mechanisms for 3D problems, and numerical methods are
 35 usually necessary.

36 Finite element limit analysis (FELA) is a powerful numerical implemen-
 37 tation that can evaluate 3D collapse loads without assuming a failure mech-
 38 anism *a priori* (Lyamin and Sloan, 2002a,b; Lyamin et al., 2007; Vicente da
 39 Silva and Antão, 2008; Krabbenhøft et al., 2008; Martin and Makrodimopou-
 40 los, 2008; Sloan, 2013). As in the conventional finite element method (FEM),
 41 FELA discretizes the domain into elements and interpolates the velocity field
 42 based on discrete values at nodes and the assumed shape functions. The opti-
 43 mal velocity field is computed by solving a large-scale optimization problem.
 44 The objective function corresponds to the limit load, and the unknown nodal
 45 velocities are constrained by enforcing kinematically admissibility. In FELA,
 46 a certain discretization of the domain (i.e., meshing) leads only to a subset
 47 of all possible velocity fields. Therefore, the limit load computed by FELA
 48 is often highly sensitive to the finite element mesh, particularly in regions of
 49 localized deformation. To maximize the solution accuracy using a minimum
 50 number of elements, adaptive mesh refinement techniques (i.e., h adaptivity)

51 have been proposed to automatically refine regions featuring large gradients
 52 (Borges et al., 1999, 2001; Lyamin et al., 2005; Martin, 2011) or large gaps be-
 53 tween upper-bound and lower-bound solutions computed on the same mesh
 54 (Ciria et al., 2008; Muñoz et al., 2009). The concept of h adaptivity has
 55 played a key role in improving the accuracy and computational efficiency of
 56 2D analyses. In contrast, 3D FELA based on adaptive mesh refinement (e.g.,
 57 Dunne and Martin, 2017) and its performance have not been investigated in
 58 great detail.

59 Another general numerical approach referred to as discontinuity layout
 60 optimization (DLO) has been developed by Smith and Gilbert (2007) and
 61 Hawksbee et al. (2013) on the basis of optimizing a velocity field consisting
 62 only of so-called velocity discontinuities, which represent infinitesimally thin
 63 zones of shearing. DLO focuses on optimizing the arrangement of these
 64 discontinuities, with the tacit assumption that the material enclosed by the
 65 discontinuities is rigid. This method searches for an optimal combination of
 66 the possible discontinuities interconnecting a *fixed* grid of nodes. Because the
 67 grid is fixed, the grid resolution has to be refined to capture intricate features
 68 or reasonably represent a continuous velocity field, which can dramatically
 69 increase the number of potential discontinuities at the cost of computational
 70 expediency (Hawksbee et al., 2013).

71 While the above-mentioned numerical approaches represent valuable tools
 72 to evaluate limit loads for 3D problems, they tend to be computationally in-
 73 tensive. In many cases, the optimal mechanism is in fact relatively simple,
 74 and the standard formulations of FELA and DLO can be unnecessarily oner-
 75 ous. Furthermore, the computational demands of existing techniques impose

76 a significant limitation for the emerging computational approach referred to
 77 in this paper as the sequential kinematic method (SKM). In SKM, kinematic
 78 solutions are sequentially computed as a means of simulating a full process of
 79 deformation, and the optimal velocity field within any particular increment
 80 is used to update the model geometry and material properties. Given its
 81 computational efficiency and stability, SKM has become a compelling alter-
 82 native to conventional techniques such as FEM for simulating problems in
 83 which capturing the evolution of material boundaries is critical (Hambleton
 84 and Drescher, 2012; Mary et al., 2013; Hambleton et al., 2014; Kong et al.,
 85 2017). In particular, SKM shows a remarkable capability in modeling the
 86 large deformation of cohesionless soils (Hambleton et al., 2014; Kashizadeh
 87 et al., 2014), which poses a significant challenge for conventional approaches.
 88 Current SKM formulations, however, are restricted to 2D. Extension to 3D
 89 has been largely halted by the lack of efficient methods to compute the op-
 90 timal velocity field within each increment of simulation.

91 In this work, we investigate the concept of r adaptivity, in combination
 92 with h adaptivity, and assess the potential of this approach for increasing
 93 computational efficiency in 3D limit analysis. Pioneered in the earlier work
 94 of Johnson (1995) and more recently explored for 2D limit analysis (Mi-
 95 lani and Lourenço, 2009; Hambleton and Sloan, 2013; Milani, 2015; He and
 96 Gilbert, 2016; Muñoz et al., 2018), r adaptivity improves the computed limit
 97 load by explicitly optimizing the *nodal positions* that control the locations
 98 of possible velocity discontinuities. Because relatively coarse meshes with
 99 suitably placed edges (velocity discontinuities) are often sufficient to obtain
 100 accurate solutions, kinematic FELA and DLO enriched with r adaptivity of-

101 fers a promising pathway for improving efficiency, as previously demonstrated
102 for 2D formulations.

103 2. Overview of the r - h adaptive approach

104 The general concept we explore in this paper is to start with a simple ve-
105 locity field, one requiring minimal computational effort, and then refine this
106 field to improve the accuracy of the computed limit load and collapse mech-
107 anism. We adopt a formulation in which the velocity field is characterized
108 by discrete regions (blocks or elements) of translational motion separated by
109 velocity discontinuities. These elements are tetrahedral by assumption, such
110 that the edges, representing velocities discontinuities, are planar. We restrict
111 our attention to material obeying the Mohr-Coulomb yield criterion and as-
112 sume that the internal friction angle ϕ and cohesion c are constant across
113 the soil mass. Similarly, the material unit weight, denoted by γ , is assumed
114 to be constant. Spatially varied ϕ , c , and γ can be included into the current
115 formulation by constructing mesh according to the soil stratigraphy, in that
116 no discontinuity spans across different layers of soils. In the case of inter-
117 layer discontinuities, the highest angle of friction and cohesion encountered
118 should be used to maintain the upper-bound status of the solution.

119 Starting from an initial arrangement (mesh) of elements, the proposed
120 r - h adaptive approach proceeds iteratively, and each iteration involves three
121 key components. First, as explained in detail in Section 3, the optimal ve-
122 locities for a fixed mesh are determined using second-order cone program-
123 ming (SOCP). Second, as described in Section 4, the nodal positions are
124 regarded as variables determined through non-linear optimization (r adap-

125 tivity). Third, elements are potentially subdivided (h adaptivity). Section
 126 5 and 6 explain this third step and how each of the three components are
 127 combined to obtain a complete solution algorithm, respectively. Section 7
 128 considers several example problems to which the algorithm is applied.

129 3. Optimization of the velocity field for a fixed mesh

130 Considering an arbitrary mesh of rigid tetrahedral elements (blocks),
 131 Hambleton and Sloan (2016) proposed a technique that utilizes second-order
 132 cone programming (SOCP) to search for a kinematically admissible velocity
 133 field that yields an optimal limit load and collapse mechanism. For com-
 134 pleteness, its mathematical formulation is summarized here.

135 A generic pair of blocks is depicted in Fig. 1(a). The velocity jump
 136 between these blocks is denoted by Δv_i and is calculated as $\Delta v_i = v_i^I - v_i^{II}$,
 137 where v_i^I and v_i^{II} are the block velocities. The superscripts I and II indicate,
 138 arbitrarily, the first and second block, and the index $i = 1, 2, 3$ indicates the

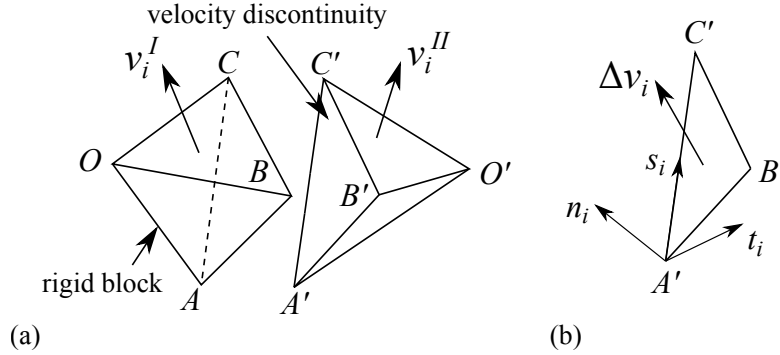


Figure 1: Schematics showing (a) 3D rigid blocks separated by a planar velocity discontinuity and (b) the definition of a local coordinate system associated with the discontinuity plane.

139 velocity component. In this work, the component associated with $i = 3$ is
 140 always in the vertical direction, and it is assumed to be positive when the
 141 velocity is upward (i.e., opposite the direction of gravity). Chen (1975) shows
 142 that for materials obeying the Mohr-Coulomb yield criterion, the energy
 143 dissipation rate along the planar velocity discontinuities between elements
 144 (blocks) can be expressed as

$$\dot{d} = cA|\Delta v_t| \quad (1)$$

145 The variable A denotes the area, and Δv_t is the tangential velocity jump with
 146 respect to the plane of the discontinuity. The absolute value is prescribed
 147 so that the dissipated power is always positive, regardless of the shearing
 148 direction. To fulfill the associative flow rule corresponding to the Mohr-
 149 Coulomb yield condition, a kinematically admissible velocity discontinuity
 150 has to meet the following “jump condition” (Chen, 1975):

$$\Delta v_n = |\Delta v_t| \tan \phi \quad (2)$$

151 The variable Δv_n denotes the normal velocity jump. By adopting the local
 152 coordinate system shown in Fig. 1(b), Eqs. (1) and (2) can be rewritten as

$$\begin{aligned} \dot{d} &= cA\sqrt{(\Delta v_i t_i)^2 + (\Delta v_i s_i)^2} \\ \Delta v_i n_i &= \tan \phi \sqrt{(\Delta v_i t_i)^2 + (\Delta v_i s_i)^2} \end{aligned} \quad (3)$$

153 Following the summation convention, the quantities $\Delta v_i t_i$ and $\Delta v_i s_i$ are dot
 154 products calculated, for example, as $\Delta v_i t_i = \Delta v_1 t_1 + \Delta v_2 t_2 + \Delta v_3 t_3$. In
 155 Eq. (3), n_i is a unit vector normal to the plane of the discontinuity, and t_i
 156 and s_i are two unit vectors parallel to the plane. These three vectors give a
 157 mutually orthogonal transformed basis for expressing the velocity vectors, as

158 depicted in Fig. 1(b). Note that in accordance with measuring the velocity
 159 jump from block I to II discussed above, the vector n_i points towards block I
 160 such that a positive normal component of the velocity jump indicates dilation.

161 In order to write Eq. (3) in a form amenable to SOCP, the quantity
 162 $\sqrt{(\Delta v_i t_i)^2 + (\Delta v_i s_i)^2}$ is replaced by a dummy variable μ :

$$\begin{aligned} \dot{d} &= cA\mu \\ \Delta v_i n_i &= \mu \tan \phi \end{aligned} \tag{4}$$

163 The dummy variable μ is then constrained as follows:

$$\mu \geq \sqrt{(\Delta v_i t_i)^2 + (\Delta v_i s_i)^2} \tag{5}$$

164 Eq. (5) is in the form of a so-called second-order cone constraint, one of
 165 the types permitted in SOCP in addition to linear equality and inequality
 166 constraints (cf. Sturm, 2002).

167 We note that the expressions given by Eq. (4) are exact only in the
 168 particular instance where strict equality is achieved in Eq. (5):

$$\mu = \sqrt{(\Delta v_i t_i)^2 + (\Delta v_i s_i)^2} \tag{6}$$

169 Equality is achieved by constructing the optimization problem such that the
 170 dummy variable μ is minimized, and thus μ is driven to equality as in Eq.
 171 (6). Application to example problems, such as those considered in Section
 172 7.3, reveals that equality is achieved in most cases. However, in the case of
 173 cohesionless material ($c = 0$) for which the dissipation $\dot{d}$ vanishes, equality
 174 is not always achieved. Nevertheless, it should be noted that, when the
 175 equality in Eq. (5) is not satisfied, the solution remains an upper bound of
 176 the true collapse load because the energy dissipation and the jump condition

are effectively computed according to a larger cohesion and friction angle,
respectively.

By equating the rate of internal energy dissipation to the work rate of
external forces for an assembly of elements (blocks), one obtains

$$\sum_{j=1}^{N_D} \dot{d}_j = - \sum_{k=1}^{N_B} \gamma V_k v_{3k} + \int_{S^*} t_i^* v_i ds + \int_S t_i v_i ds \quad (7)$$

In Eq. (7), N_D and N_B are the number of discontinuity planes and the number
of blocks, respectively, and subscripts j and k are used to indicate quantities
corresponding to the j^{th} discontinuity plane and the k^{th} block. The variable
 V_k denotes the volume of the k^{th} block, a readily computed constant for a
fixed mesh. The three terms on the right side of Eq. (7) represent the work
rate of body forces, fixed surface tractions (t_i^*) and tractions along the surface
where the limit loads is evaluated (t_i), respectively.

Drescher (1991), Sloan (1995), and Michalowski (2001) among others
show how Eq. (7) can be manipulated to obtain various expressions of the
limit load. A case encompassing all examples considered in Section 7 is
that the direction of t_i is fixed, or known *a priori*, and velocities along the
boundary S are uniform, as could occur for a rigid footing or translational
retaining wall. In this instance, the unknown traction t_i is expressed in
terms of a fixed traction t_i^* as $t_i = \lambda t_i^*$, where $\lambda \geq 0$ is an unknown multi-
plier dictating the magnitude of the limit load. The last term in Eq. (7) is
 $\int_S t_i v_i ds = v_i \int_S \lambda t_i^* ds = \lambda v_i F_i^*$, where $F_i^* = \int_S t_i^* ds$ is the resultant force.
The magnitude of the velocity is arbitrary (cf. Chen, 1975), and thus one
can write $v_i F_i^* = \alpha$, where α is an arbitrary constant. Equation (7) can then

199 be manipulated to write:

$$\lambda = \frac{1}{\alpha} \left(\sum_{j=1}^{N_D} \dot{d}_j + \sum_{k=1}^{N_B} \gamma V_k v_{3k} - \int_{S^*} t_i^* v_i ds \right) \quad (8)$$

200 Here we assume α is unity (a value of 1 with appropriate units) for convenience. Depending on the distribution of the fixed tractions t_i^* , the final term in parenthesis in Eq. (7) can be integrated to obtain a sum over the unknown velocities, viz.

$$\int_{S^*} t_i^* v_i ds = \sum_{l=1}^{N_F} \beta_{il} v_{il} \quad (9)$$

204 In Eq. (9), N_F is the number of elements with fixed tractions, and β_{il} ($i = 1, 2, 3; l = 1, \dots, N_F$) are constant coefficients. The notation v_{il} again indicates the i^{th} velocity component of the l^{th} element.

207 Finally, the optimization of the velocity field for a fixed mesh is written
208 in the standard form of SOCP as follows:

$$\begin{aligned} \min \quad & \lambda = \sum_{j=1}^{N_D} \dot{d}_j + \sum_{k=1}^{N_B} \gamma V_k v_{3k} - \sum_{l=1}^{N_F} \beta_{il} v_{il} \\ \text{s.t.} \quad & \Delta v_{ij} n_{ij} = \mu_j \tan \phi \quad j = 1, \dots, N_D \\ & \dot{d}_j = c A_j \mu_j \quad j = 1, \dots, N_D \\ & \mu_j \geq \sqrt{(\Delta v_{ij} t_{ij})^2 + (\Delta v_{ij} s_{ij})^2} \quad j = 1, \dots, N_D \end{aligned} \quad (10)$$

209 For a load resisting collapse, where the work rate of the unknown tractions
210 on the velocity is negative, the kinematic theorem of limit analysis leads to
211 a lower bound on the true collapse load (cf. Drescher, 1991). To compute
212 such a lower bound, Eq. (10) is converted to a maximization problem by
213 minimizing the negative of the objective function. In this work, the Mosek

214 toolbox integrated with MATLAB (Mosek, 2015) is employed to solve the
215 SOCP problem.

216 Upon solving the SOCP problem of Eq. (10), one obtains an optimal
217 value for the load multiplier, denoted by λ_{opt} . The computed bound on the
218 true collapse load is then simply

$$F_i = \lambda_{opt} F_i^* \quad (11)$$

219 4. Optimization of nodal positions (r adaptivity)

220 The bound on the limit load computed using Eq. (10) depends strongly
221 on the positions of the nodes within the mesh that define the locations of
222 potential velocity discontinuities. In particular, the optimal velocity field
223 and load multiplier λ_{opt} depend on the coordinates of the nodes that are not
224 constrained by boundary conditions or symmetry, and are therefore free to
225 move. The coordinates of these nodes are denoted by x_{im} . Index i again
226 gives the component ($i = 1, 2, 3$), and index m ($m = 1, \dots, N_R$) identifies each
227 of the free nodes.

228 For the purpose of optimizing the nodal positions, a non-linear optimiza-
229 tion problem is formulated as follows:

$$\begin{aligned} \min \quad & \lambda_{opt}(x_{im}) \\ \text{s.t.} \quad & V_k(x_{im}) \geq 0 \quad k = 1, \dots, N_B \\ & x_{im}^l \leq x_{im} \leq x_{im}^u \end{aligned} \quad (12)$$

230 This non-linear optimization is nested with the SOCP described above, in
231 that the objective function in Eq. (12) is the load multiplier computed for
232 a given set of nodal positions x_{im} ($i = 1, 2, 3; m = 1, \dots, N_R$), defined and

233 evaluated in precisely the same way as in Section 3. To prevent the inter-
 234 penetration of elements and ensure computational stability, the first set of
 235 constraints in Eq. (12) requires that element volume V_k ($k = 1, \dots, N_B$) is
 236 always positive. It should be noted that we permit the possibility $V_k = 0$, thus
 237 allowing elements to collapse to transition layers with zero thickness. The
 238 variables x_{im}^l and x_{im}^u appearing in the second set of inequality constraints
 239 define allowable limits for certain nodal position components. For instance,
 240 the z -coordinate of the ground surface is an upper bound on the position of
 241 all nodes along the z -direction.

242 Due to boundary conditions and symmetry, some of the position compo-
 243 nents (x , y , and z) are fixed. Rather than imposing constraints, the total
 244 number of free variables introduced in the non-linear optimization problem
 245 of Eq. (12) is condensed from $3N_R$ to DOF , where $DOF = 3N_R - N_{FC}$ and
 246 N_{FC} is the total number of fixed position components.

247 As the objective function and constraints are non-linear functions of the
 248 free (unknown) variables x_{im} , the optimization problem of Eq. (12) falls
 249 within the general domain of non-linear constrained optimization. A pre-
 250 liminary study employs two algorithms embedded in the FMINCON solver
 251 of MATLAB to solve this problem: the interior point method (IPM) and
 252 sequential quadratic programming (SQP). Both methods represent the state
 253 of the art in solving general constrained optimization problems. It is found
 254 that these two methods can achieve similar solutions. However, IPM requires
 255 more iterations, and during some iteration processes it diverges (i.e., the ob-
 256 jective function increases rather than decreases). Accordingly, SQP is used
 257 throughout this work. It should be noted that the theoretical reason why

258 SQP outperforms IPM remains unclear. This is due in part to the lack of an
259 explicit expression for the objective function in the constrained optimization
260 problem (i.e., the objective function itself is the SOCP problem defined in
261 Eq. (10)).

262 To determine when to stop the iterations for solving the optimization
263 problem of Eq. (12), we adopt two criteria, and the satisfaction of either
264 one is assumed to signal the convergence to a solution. Specifically, the
265 optimization ends once (1) the quantity referred to as “first-order optimality”
266 is lower than a tolerance, opt_{tol} , or (2) the norm of the vector containing the
267 changes of nodal positions during an iteration is lower than a tolerance,
268 Δx_{tol} . First-order optimality, described in greater detail by Nocedal and
269 Wright (2006), is a well-known and widely used measure of how close the
270 current solution is to optimal. We use the second criterion to cease iterations
271 when r adaptivity produces only minor perturbations that lead to marginal
272 improvement the computed limit load. The following tolerance values are
273 employed in this work: $\text{opt}_{tol} = 1E^{-2}$ and $\Delta x_{tol} = 1E^{-2}$. The usage of
274 lower tolerances increases the number of iterations but does not noticeably
275 improve the solution. For detailed descriptions of the above stopping criteria
276 and their implementation in MATLAB, the reader is referred to Nocedal and
277 Wright (2006) and The MathWorks, Inc (2018).

278 5. Element subdivision (h adaptivity)

279 Once r adaptivity is applied to optimize the limit load and velocity field
280 for a particular mesh topology (element number and connectivity), further
281 improvement of the solution requires either uniformly or selectively refining

282 the mesh. This section proposes a strategy to refine the mesh by selectively
283 dividing elements, such that refinement will only be performed as needed and
284 at locations that potentially improve the solution.

285 Any subdivision strategy must decide where to refine the mesh based
286 on certain *a posteriori* indicators (i.e., information derived from the current
287 computation). For a rigid block system, a simple indicator is the magni-
288 tude of the velocity jump, which is proportional to the integral of strain rate
289 over the infinitesimally thin layers between adjacent elements (Chen, 1975)
290 represented as velocity discontinuities. The magnitude of the velocity jump
291 therefore identifies regions characterized by high strain rate, and mesh re-
292 finement in these regions typically has the highest potential for improving
293 the solution. This concept is similar to the adaptive mesh refinement pro-
294 posed by Martin (2011) for FELA, which attempts to evenly distribute the
295 integral of the maximum shear strain rate over all elements, such that the
296 concentration of elements reflects the intensity of the shearing rate (change
297 of velocity). The specific subdivision criterion postulated in this work is to
298 subdivide elements sharing an edge for which the magnitude of the velocity
299 jump is greater than a tolerance Δv_{tol} , i.e., $\sqrt{\Delta v_i \Delta v_j} \geq \Delta v_{tol}$.

300 The flow chart within the dashed box of Fig. 2 presents the basic algo-
301 rithm iterated over all elements to perform the subdivision. This algorithm
302 first filters out elements with nearly zero velocity or small volume through
303 prescribed tolerances v_{tol} and V_{tol} , respectively. The former filtering prevents
304 unnecessary refinement in stationary regions, and the latter contributes to
305 forming the best overall shape of the mechanism, excluding the partition of
306 small elements that tend only to result in small and localized improvement

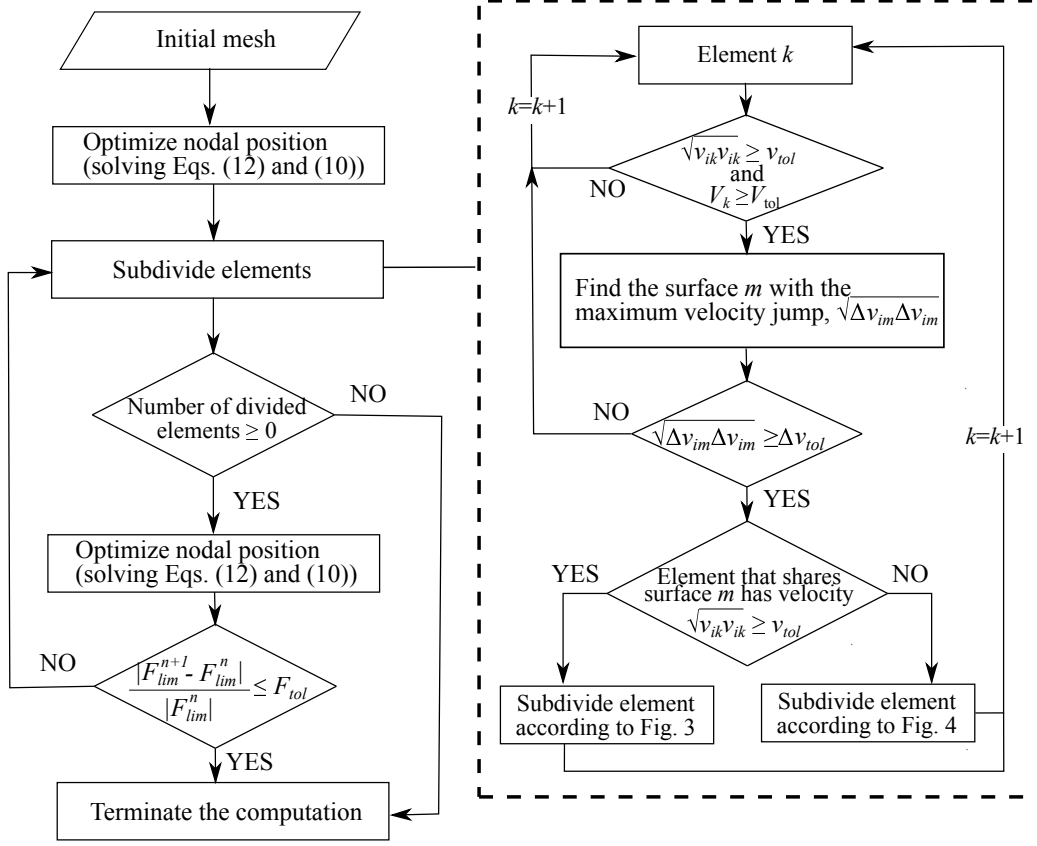


Figure 2: Computation flow chart of the proposed r - h adaptive approach.

of the collapse mechanism. Each element that passes this first screening and has edges with $\sqrt{\Delta v_i \Delta v_i} \geq \Delta v_{tol}$ will be subdivided according to either Fig. 3 or Fig. 4, depending on whether this velocity jump is between two moving elements (i.e., both have velocity greater than v_{tol}), or between a moving element and a stationary region.

As depicted in Fig. 3, when the targeted velocity jump is between two moving elements, we propose two different approaches to subdivide the element corresponding to the subfigures (a) and (b). The adoption of one

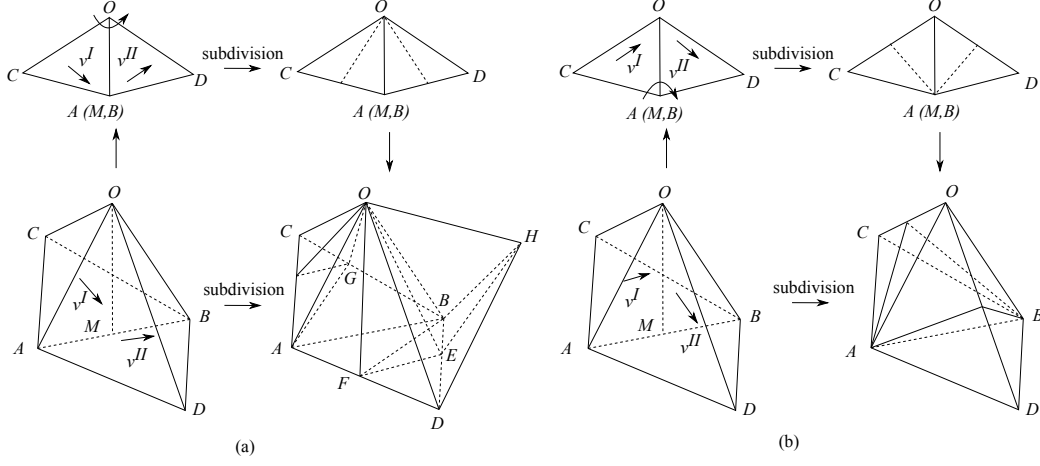


Figure 3: Schematic showing the subdivision of a pair of blocks based on the flow direction of the local velocity field.

of these two alternatives depends on *the flow direction of the local velocity field*. Fig. 3 shows the two possibilities: a flow tending to “rotate” about the point O , as shown in subfigure (a), or rotate about the axis AB , as shown in subfigure (b). Identifying this flow direction is important because regional velocity jumps can be reduced (smoothed) when more elements are added aligning with this direction.

For the scenario shown in Fig. 3(a), the blocks $OABC$ and $OADB$ are divided so that the newly added discontinuities radiate from the point O and bisect the edges AC , BC , AD and BD . Note that for illustration purposes, we have assumed the surface OAB possesses the maximum velocity jump for both blocks; otherwise, only one block is subdivided. Fig. 3(a) also shows that the subdivision of the targeted element $OADB$ adds a new node E to the edge BD , which is shared by an adjacent element $ODBH$. These neighboring elements will automatically be partitioned by new discontinuities

329 passing through the new nodes (e.g., the new discontinuity OEH is created
 330 to pass through the node E in Fig. 3(a)). Without such partition of neigh-
 331 boring elements, subsequently changing the positions of the new nodes (e.g.,
 332 the node E) can lead to interpenetration or gaps between the newly formed
 333 elements (e.g., the blocks $OFEB$ and $OFDE$) and those that already ex-
 334 isted (e.g., the block $ODBH$). Moreover, subdividing these adjacent blocks
 335 ensures that the newly formed discontinuities are connected (e.g., the dis-
 336 continuities OFE and OEH), thus enabling immediate benefits from the r
 337 adaptivity. Due to the fact that only tetrahedral elements are considered,
 338 some secondary discontinuities (e.g., the discontinuities OAG and OFB in
 339 Fig. 3(a)) are added during the subdivision process. Extending the proposed
 340 approach to other element shapes would eliminate this requirement.

341 When the local velocity field features the characteristics shown in Fig. 3(b),
 342 the newly added discontinuities radiate from the axis AB and bisect the edges
 343 OC and OD , and there are many possible ways to distinguish the above two
 344 different flow directions. The one employed in this work is given by

$$\begin{cases} \text{rotate about } AB & \text{if } \Delta v_i r_i < 0 \\ \text{rotate about } O & \text{if } \Delta v_i r_i \geq 0 \end{cases} \quad (13)$$

345 where $\Delta v_i = v_i^I - v_i^{II}$, with v_i^I and v_i^{II} denoting the element velocities pointing
 346 toward and away from the shared surface OAB shown in Fig. 3, respectively.
 347 The variable r_i in Eq. (13) represents a unit vector pointing from O to M . It
 348 is used as a reference direction for distinguishing the direction of the velocity
 349 jump. When a pair of elements have velocities that both point toward or
 350 away from the interface (OAB in Fig. 3), they will not be divided in the
 351 current iteration, due to the ambiguity of the flow direction.

Figure 3 shows only one of three possible permutations, namely the flow direction of the local velocity field can also rotate about the other two pairs: (1) the point A paired with the axis BO and (2) the point B paired with the axis AO . These three possibilities are distinguished by projecting the velocity jump to the three edges of the triangle OAB . The edge with the least projection is the one to which the velocity jump has the greatest perpendicular component, and thereby the one about which the local velocity flow tends to rotate. Mathematically, this criterion can be expressed as

$$\begin{cases} \text{rotate about } O/AB & \text{if } |\Delta v_i o_i| \leq \min(|\Delta v_i p_i|, |\Delta v_i q_i|) \\ \text{rotate about } A/BO & \text{if } |\Delta v_i p_i| \leq \min(|\Delta v_i o_i|, |\Delta v_i q_i|) \\ \text{rotate about } B/AO & \text{if } |\Delta v_i q_i| \leq \min(|\Delta v_i o_i|, |\Delta v_i p_i|) \end{cases} \quad (14)$$

where o_i , p_i and q_i denote vectors along edges AB , BO and AO , respectively.

Elements adjacent to stationary regions are subdivided as illustrated in Fig. 4. Specifically, the element is divided by creating a new discontinuity that radiates from the point O and bisects the edges AC and BC . The

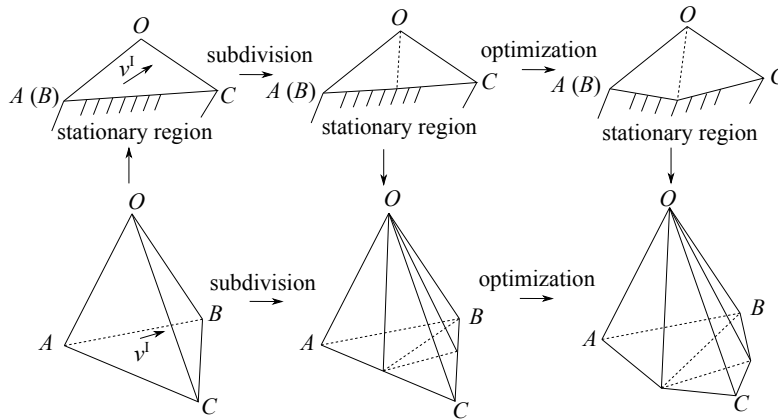


Figure 4: Schematic showing subdivision of a moving block adjacent to stationary region.

364 decision as to which edges to bisect are determined in a manner similar to
 365 Eq. (14). In the rightmost figures, we show the new nodes in their optimized
 366 positions (off of plane ABC) to illustrate that this type of subdivision enables
 367 an accurate resolution of the boundary between moving material and the
 368 stationary region, which is typically a discontinuity whose shape is not known
 369 beforehand.

370 6. Algorithm summary

371 The complete algorithm for the proposed r - h adaptive method is summa-
 372 rized in the main flow chart of Fig. 2. The computations start by optimizing
 373 the nodal positions of the initial mesh. Then, the algorithm repeats the cycle
 374 of subdividing elements and adjusting nodal positions, until satisfying either
 375 of the following two criteria: (1) the relative improvement of the limit load
 376 between two consecutive subdivisions is less than a prescribed tolerance, de-
 377 noted by F_{tol} , or (2) no element needs to be subdivided. It should be noted
 378 that, in the above-mentioned cycle, any h adaptivity step is immediately fol-
 379 lowed by an r adaptivity step. The reason why we do not allow consecutive
 380 h adaptivity steps will be elaborated by the numerical examples detailed in
 381 Section 7.

382 As in any numerical approach, the question arises as to how to select the
 383 various tolerances introduced above. For the numerical examples discussed
 384 later, trial and error revealed that the following choices of tolerances give
 385 satisfactory performance: $F_{tol} = 0.1$, $\Delta v_{tol} = v_0$, $v_{tol} = 0.01v_0$, where v_0 de-
 386 notes the magnitude of the velocity along the boundary where the limit load
 387 is evaluated. Because the volume filtering mechanism described above can

388 potentially stop subdivision prematurely, a small value of $1E^{-3}b^3$ is assigned
389 to the tolerance V_{tol} , where b is the largest dimension of the loading area.

390 7. Example problems

391 To explore the performance of the proposed method, three examples are
392 studied: (1) bearing capacity of a square foundation on cohesionless soil or
393 purely cohesive soil; (2) passive uplift resistance of a square, horizontal anchor
394 embedded in cohesionless soil; and (3) passive resistance of a rectangular
395 retaining wall in cohesionless soil.

396 7.1. Bearing capacity of a square foundation

397 The limit load for a square surface foundation of width b on cohesionless
398 soil can be expressed as

$$F = \frac{1}{2}\gamma b^3 N_{\gamma s} \quad (15)$$

399 In Eq. (15), the dimensionless quantity $N_{\gamma s}$ is a function of the internal fric-
400 tion angle ϕ and the interfacial roughness between the footing and the soil.
401 The subscript “s” is used to distinguish this factor, for a square foundation,
402 from the 2D (plane strain) bearing capacity factor commonly denoted as
403 N_γ . Exact values for N_γ were obtained numerically by Martin (2005), who
404 performed detailed calculations based on the method of characteristics and
405 utilized, notably, adaptive subdivision in his approach. In 3D, exact theoret-
406 ical solutions remain elusive, and $N_{\gamma s}$ in particular is an unknown function.
407 However, upper bounds obtained through limit analysis have been evaluated
408 semi-analytically and numerically (Michalowski, 2001; Krabbenhøft et al.,
409 2008; Lyamin et al., 2007). This work models cohesionless soils by assigning

410 zero-valued cohesion c to the dissipated power (Eq. 10). The unit weight of
 411 the soil γ and the footing width b are each assumed to be 1 for ease in inter-
 412 preting $N_{\gamma s}$. The relative slip between the footing and the soil is prevented
 413 (i.e., perfectly rough) in the simulation.

414 To initiate the computation, one has to guess an initial mesh. For refer-
 415 ence, we consider the mechanism constructed by Michalowski (2001) rendered
 416 in Fig. 5(a). This mechanism is characterized by a single pyramidal block
 417 that moves downward vertically with the foundation and four adjacent re-
 418 gions composed of rigid blocks truncated by conical surfaces. For clarity, Fig.
 419 5 shows only one of the four regions. By comparison, the starting guess con-
 420 sidered in this work is extremely simple. It is depicted in Fig. 5(b). Taking
 421 advantage of the four-fold symmetry (i.e., OMN shown in Fig. 5(a) represents
 422 a 45° slice of the footing), the mesh consists of only three elements (blocks),
 423 one directly beneath the foundation and two that are adjacent. Initially, one

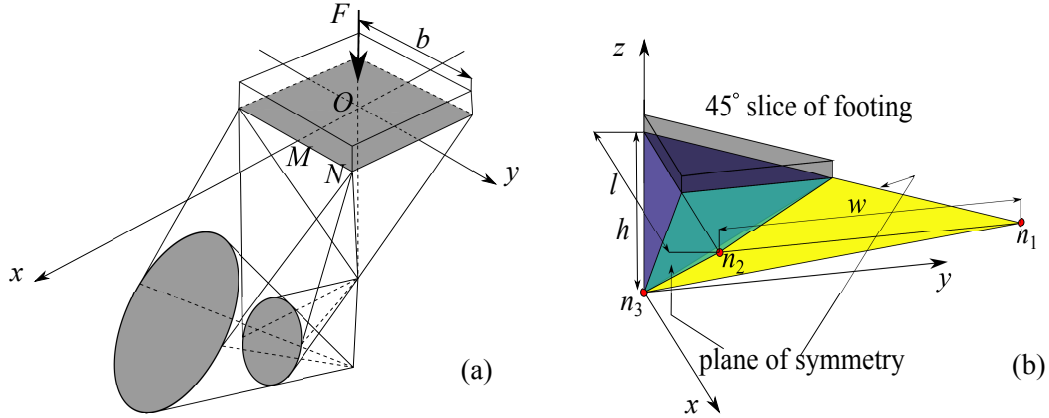


Figure 5: Bearing capacity of a rough rigid square foundation on cohesionless soils: (a) multi-block mechanism (adapted from Michalowski (2001)); (b) initial mesh assumed in the r - h adaptive approach.

424 has to guess the positions of the nodes n_1 , n_2 , and n_3 (or equivalently the
425 values of three geometric variables h , l and w in Fig. 5(b)). Throughout the
426 r - h adaptive optimization procedure (Fig. 2), these nodes are constrained
427 to move parallel to the plane of symmetry in which they reside, $y = 0$ or
428 $x = y$, as are any nodes within these planes added through adaptive sub-
429 division. Additionally, the components of velocity normal to the planes are
430 constrained to be zero.

431 When the friction angle is high, the jump condition given by Eq. (2)
432 becomes increasingly restrictive with respect to finding a kinematically ad-
433 missible velocity field for a particular mesh. Consequently, the existence of
434 a feasible solution for SOCP becomes sensitive to the mesh geometry, and
435 selecting initial values for the above geometric variables becomes challeng-
436 ing. This issue was resolved by sequentially optimizing the nodal positions
437 while gradually increasing the friction angle. In other words, one can start
438 the computation by (1) introducing a low, fictitious friction angle denoted
439 by ϕ_0 , (2) optimizing the nodal positions, and (3) using the optimized mesh
440 as a starting guess to obtain a feasible initial solution for a higher friction
441 angle. The procedure is repeated until the true friction angle is reached. In
442 this work, the starting guess in all cases was $h = b/2$, $w = b$, and $l = b$ with
443 $\phi_0 = 10^\circ$.

444 The solid line in Fig. 6 shows the computed values of N_{γ_s} for $\phi = 35^\circ$
445 as they vary for each iteration of the SQP algorithm utilized within the
446 proposed r - h adaptive solution procedure to solve Eq. (12). The figure shows
447 that the computed upper bound on N_{γ_s} rapidly decreases as the iteration
448 number increases, highlighting the sensitivity of the solution to the mesh,

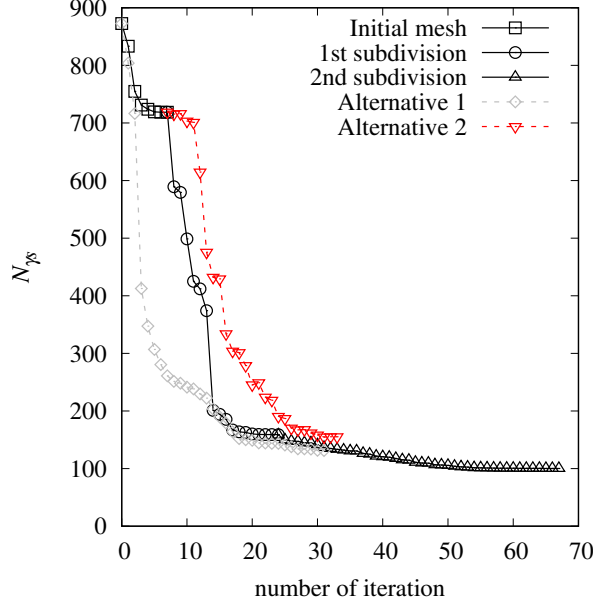


Figure 6: Variation of the N_{γ_s} value as a function of the iteration numbers in the non-linear optimization ($\phi = 35^\circ$).

449 and thus also revealing the effectiveness of r adaptivity. In this example,
 450 the method resulted in two subdivisions. The initial mesh and the meshes
 451 corresponding to these subdivisions are presented in Figs. 7(a)-(c), wherein
 452 the number of elements after each subdivision is also provided. Prior to
 453 each subdivision, the convergence curve becomes flat, signaling that better
 454 upper bounds cannot be reached for the current mesh. Through the use of
 455 h adaptivity, the computed limit load can be further reduced, and a faster
 456 convergence rate can be recovered (e.g., 1st subdivision in Fig. 6). The
 457 reason why h adaptivity is effective is revealed in Fig. 7. Comparing the
 458 initial mesh to those obtained after subdivision, the approach enables the
 459 creation of more velocity discontinuities radiating outward from the edge

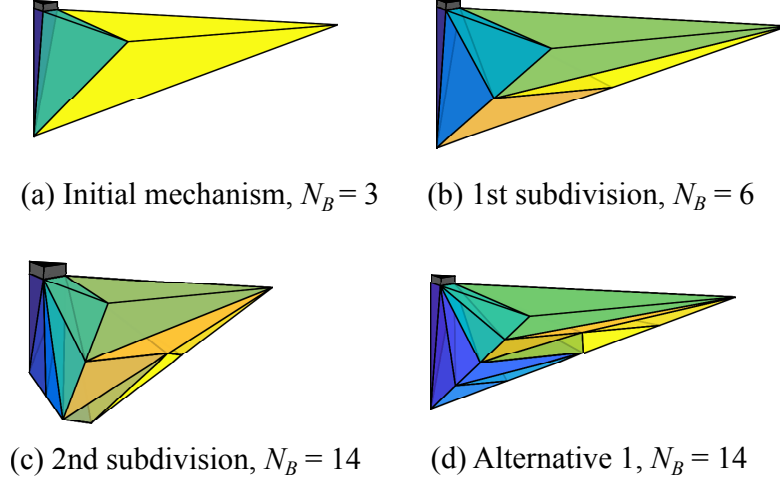


Figure 7: Mesh before and after element subdivisions ($\phi = 35^\circ$).

of the footing. Moreover, the lowermost part of the collapse mechanism is gradually divided in a way that the above radial discontinuities can extend all the way to the boundary of the region of failing (moving) soil. Both features are important in forming a radial shearing zone, which accommodates the rotation of the principal directions of strain.

Figure 6 highlights the fact that the rate of improvement in the solution generally diminishes as the r - h adaptive iterations proceed. This response may be attributed to two possible explanations. The first hypothesis is that the smaller element sizes obtained through h adaptivity constrain the magnitude of nodal position changes that can occur in the optimization, due to the imposed non-linear constraints requiring no interpenetration between elements (see Eq. (12)). This reduces the amount that nodes are able to perturb around their current positions, thereby demanding more r -adaptive iterations to achieve a better mechanism. The second hypothesis is that, as the current mechanism is closer to the optimum, the optimization algorithm

475 adopts smaller step size (i.e., nodal perturbation during one iteration), and
476 consequently the rate of improvement is reduced.

477 To test these two hypotheses, we consider an alternative initial mesh with
478 the same overall geometry as the original one (Fig. 7(a)) but with the same
479 connectivity and number of elements as in the final solution (Fig. 7(c)). Com-
480 pared with the original starting guess, this new initial mechanism (Fig. 7(d))
481 simply has smaller initial element sizes. The dashed line in Fig. 6 designated
482 by “Alternative 1” shows that the r adaptive iterations are more effective
483 for the new initial mesh with smaller element sizes. This reveals that the
484 deterioration in the effectiveness of r adaptivity is not related to the number
485 and size of elements but rather due to the fact that an optimal mechanism is
486 approached (the second of the two hypotheses above). One might infer from
487 the above discussion that starting from a more refined mesh is generally
488 more effective, given that better results are achieved with fewer iterations.
489 However, this is not the case, since the refined solution with element edges
490 (velocity discontinuities) placed at strategic locations is known only after
491 refinements are obtained through iterations of r - h adaptivity.

492 We use the data corresponding to “Alternative 2” in Fig. 6 to illustrate
493 why an r adaptive step is employed immediately following any h adaptive
494 step, as described in Section 6. In Alternative 2, two consecutive h adaptive
495 steps are performed on the initial mesh depicted in Fig. 7(a). The nodal
496 positions of this refined mesh are then optimized using r adaptivity. Fig-
497 ure 6 shows that after multiple h adaptive steps, r adaptivity becomes less
498 efficient compared with the proposed algorithm. This can be explained as
499 follows. When only subdividing elements without optimizing nodal positions,

500 the elements on both sides of the new discontinuities have the same veloc-
 501 ity as if the original elements have yet to be subdivided, and there are no
 502 velocity jumps across these new discontinuities. These new discontinuities
 503 with zero velocity jumps do not provide effective information regarding how
 504 to refine the mesh (see Section 5). This analysis shows that simply increas-
 505 ing the number of elements often does not lead to an improved solution. It
 506 underscores the merit of the proposed approach, which starts from a simple
 507 mesh that is progressively refined through combined r - h adaptivity.

508 Figure 8 presents the N_{γ_s} values computed with the proposed method for
 509 friction angles from $\phi = 15^\circ$ to 35° . For the purpose of comparison, the exist-
 510 ing semi-analytical solution of Michalowski (2001) and the FELA results of

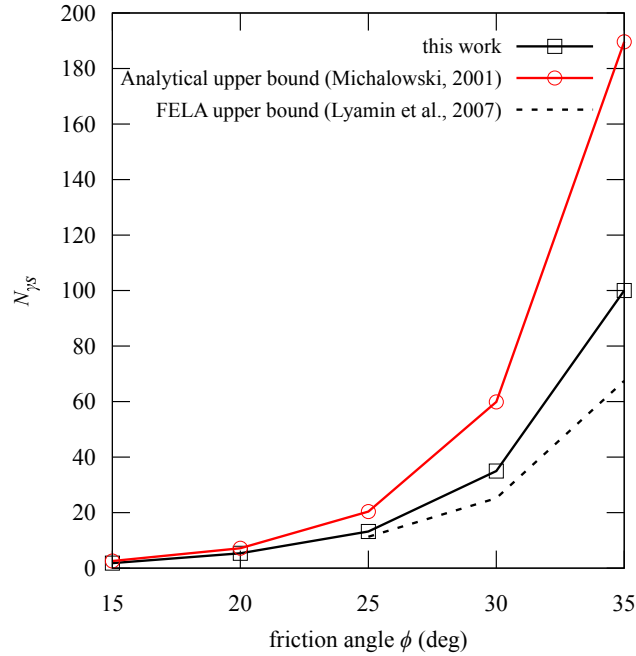


Figure 8: Comparison of N_{γ_s} values computed from the proposed method and existing solutions.

511 Lyamin *et al.* (2007) are included in the figure. Figure 8 shows that the pro-
 512 posed method gives a better (smaller) solution than the analytical approach,
 513 with the improvement increasing as the friction angle grows. On the other
 514 hand, the $N_{\gamma s}$ values computed in the present study are larger than the ones
 515 given by FELA, with the difference again tending to increase as the friction
 516 angle increases. Such a discrepancy between these two methods might be at-
 517 tributed to the continuous deformation allowed within elements in FELA.
 518 Through the use of rigid elements, the implementation presented in this work
 519 is potentially restrictive in the manner in which it accommodates the dilation
 520 of soils with large friction angles. Nevertheless, the ability of such a simple
 521 approach, and relatively simple collapse mechanism, to capture reasonable
 522 values of the limit load for such a challenging problem is remarkable.

523 Figure 9 compares the computed failure mechanisms for different friction

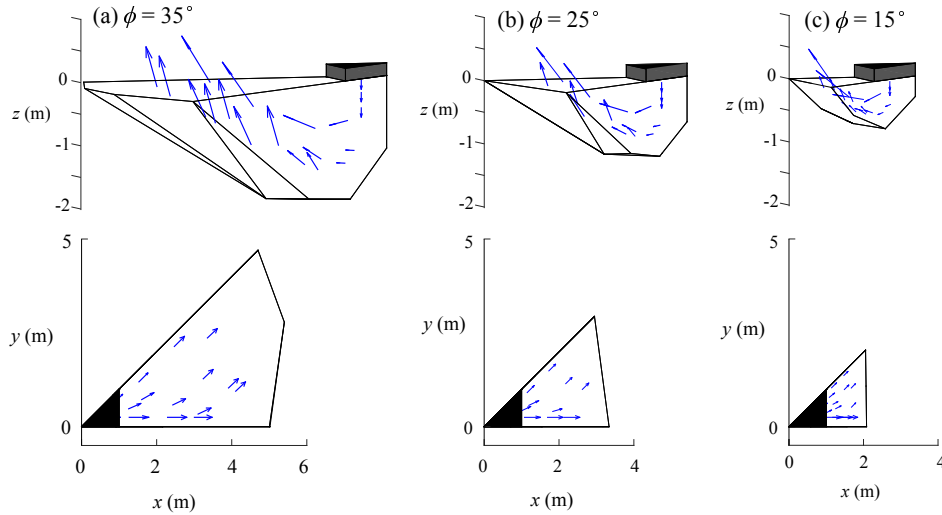


Figure 9: Optimal failure mechanism beneath the square foundation computed with the proposed approach: (a) $\phi = 35^\circ$; (b) $\phi = 25^\circ$; (c) $\phi = 15^\circ$.

524 angles. As a matter of clarity, some block edges are removed from the figure.
 525 It should be emphasized that while all three mechanisms start from the
 526 same guess in terms of the initial mesh (Fig. 5(b)), they automatically evolve
 527 depending on the friction angle of the material. For larger friction angles, the
 528 failure mechanism extends both horizontally and vertically a larger distance
 529 compared to solutions with lower friction angles.

530 To further test the proposed method, especially against existing tech-
 531 niques, the problem of a square foundation on cohesive soil is analyzed. The
 532 limit load for this problem, first considered by Shield and Drucker (1953),
 533 can be expressed in terms of the soil cohesion c as

$$F = cb^2N_{cs} \quad (16)$$

534 where N_{cs} is a constant. Computations were completed in the same manner
 535 as for $N_{\gamma s}$, using the same initial guess for the mesh as described above
 536 (Fig. 5(b); $h = b/2$, $w = b$, and $l = b$). The unit weight of soils is assumed
 537 to be zero and the cohesion c is equal to 1.

538 Table 1 compares the N_{cs} values computed in this study to those ob-
 539 tained semi-analytically (Michalowski, 2001), using FELA (Vicente da Silva
 540 and Antão, 2008), and using DLO (Hawksbee et al., 2013). The recorded
 541 or reported computation time for each numerical method is also included.
 542 Calculations in this study were completed on a PC equipped with an Intel
 543 i7-4790 processor (3.6 GHz; 4 cores) and 8 GB memory. Results from FELA
 544 were obtained by distributing computations over 5 or 18 PCs, where each
 545 PC was equipped with a single core processor clocked at 3.0 GHz (Intel Pen-
 546 tium IV) and 512 MB memory. The DLO computations were performed on
 547 a workstation equipped with an AMD Opteron 6140 processor (2.6 GHz; 8

Table 1: Comparison of computed N_{cs} values by different methods with corresponding computation costs (wall-clock time: the total processing time, including time spent on pre-processing, kernel computation through the FMINCON function and MOSEK, and post-processing; MOSEK time: the processing time spent on solving the second-order cone programming through MOSEK).

Analytical		FELA		DLO		This work			
upper bound		upper bound							
N_{cs}	N_{cs}	wall-clock times (s)	node spacing	N_{cs}	MOSEK times (s)	subdivision	N_{cs}	wall-clock times (s)	MOSEK times (s)
6.56	6.05	2000	1/2	6.52	0.02	0	8.27	0.96	0.2
		to	1/4	6.41	13	1	6.67	2.0	0.4
		15000	1/6	6.22	6400	2	6.44	4.4	0.9

548 cores) and 8 GB memory, and only the CPU times for SOCP with Mosek
549 were reported. Because the basic computation unit in the proposed r - h adap-
550 tive approach is solving Eq. (10) using SOCP, the CPU times for executing
551 Mosek are separated from the total wall-clock times. Due to differences in
552 the hardware and the particulars of programming (e.g., language and code
553 optimization), the computation times reported in Table 1 are merely indi-
554 cators of the computation cost rather than strictly comparable performance
555 measures.

556 Table 1 shows that DLO and the r - h adaptive approach provide reason-
557 ably accurate estimates of the limit load (better than the analytical solution),
558 and that FELA gives the least upper bound (best estimate of the limit load).

559 Compared to DLO and FELA, for which accuracy and computation time de-
560 pend strongly on the element size or grid spacing, the r - h adaptive approach
561 displays a significant improvement in the computed limit load without an
562 exorbitant increase in computational cost. This difference can be attributed
563 to the fact that uniform mesh or grid refinement tends to add a large num-
564 ber of additional unknowns that do not contribute towards improving the
565 solution. Finally, we note that the CPU times for running Mosek in this
566 work are only a small portion of the total times, thus suggesting that the
567 reported computational times can be potentially reduced by utilizing more
568 efficient optimization schemes and programming languages for the non-linear
569 optimization problem posed by r adaptivity (Section 4).

570 7.2. Uplift resistance of a plate anchor in cohesionless soil

571 With reference to the collapse mechanism considered by Murray and Ged-
572 des (1987), Fig. 10(a) illustrates the problem of a horizontal anchor problem
573 embedded at depth h . The anchor is square with sides of length, b , and the
574 material is assumed to be cohesionless. The status of “immediate breakaway”
575 is considered, which implies that the underside of the anchor loses contact
576 with the soil. The ultimate uplift force F is expressed as

$$F = \gamma h b^2 N_{\gamma b} \quad (17)$$

577 The factor $N_{\gamma b}$ is referred to as the anchor break-out factor, and its value
578 depends on ϕ , the ratio of the embedment depth to the anchor width (h/b),
579 and friction at the soil-anchor interface. For a fixed friction angle of $\phi = 30^\circ$,
580 this example considers the $N_{\gamma b}$ factors corresponding to varying values of
581 h/b . Here, as in the previous example, we assume zero-valued cohesion to

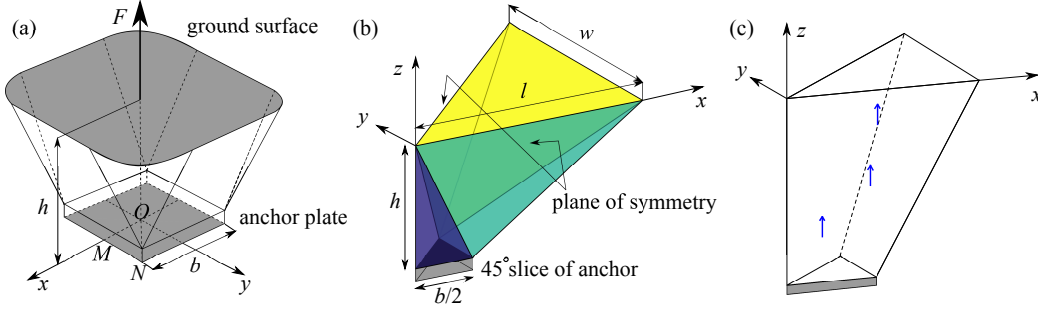


Figure 10: Uplift of an anchor in cohesionless soil: (a) collapse mechanism considered by Murray and Geddes (1987); (b) initial mesh used in the r - h adaptive approach; (c) typical collapse mechanism computed with the r - h adaptive approach.

582 model cohesionless soils. The unit weight of the soil γ and the anchor width
 583 b are each assumed to be 1. Above the anchor, a perfectly rough interface
 584 is simulated by eliminating relative movement between the anchor and the
 585 soil.

586 Fig. 10(b) depicts the initial mesh selected for the r - h adaptive approach.
 587 As in the previous example, symmetry is invoked to reduce the model to a
 588 45° slice of the anchor (i.e., slice OMN in Fig. 10(a)), where the planes $y = 0$
 589 and $x = y$ represent the planes of symmetry. The initial mesh is again one
 590 of the simplest conceivable, and it consists of three elements. The geometric
 591 variables l and w are initially assumed to be $2h$, which leads to feasible initial
 592 solutions for all embedment ratios.

593 Figure 11 compares the $N_{\gamma b}$ values computed with the r - h adaptive ap-
 594 proach to values obtained in previous works: those obtained with the ana-
 595 lytical solution of Murray and Geddes (1987) and the 3D DLO analysis of
 596 Hawksbee et al. (2013). Satisfactory agreement between the three methods
 597 can be observed for both shallow and deep embedment. Figure 10(c) presents

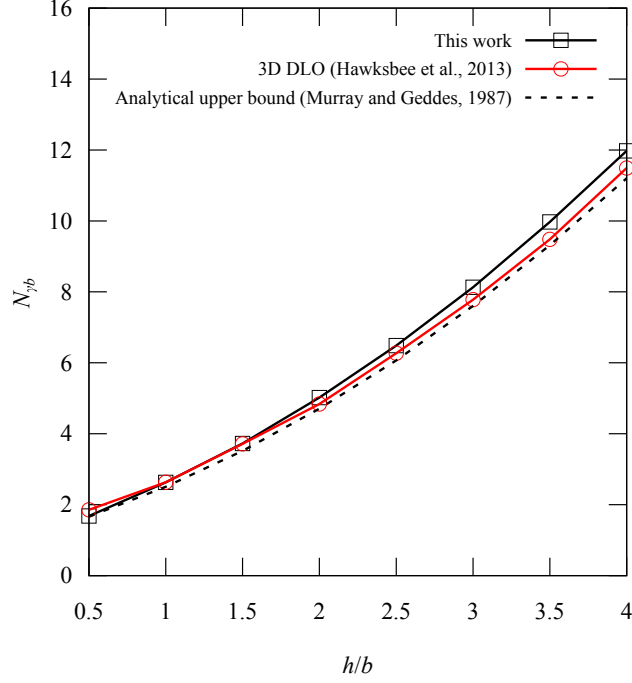


Figure 11: $N_{\gamma b}$ values computed with the r - h adaptive approach compared to existing solutions.

598 a typical collapse mechanism computed by this work. The mechanism is char-
 599 acterized by single active velocity discontinuity that extends from the edge of
 600 the anchor to the ground surface and bounds a plug of material that moves
 601 upward with the anchor. This mechanism is similar to the one constructed
 602 by Murray and Geddes (1987) (Fig. 10(a)), but it differs with respect to the
 603 the conical surfaces assumed at the edges of the collapse mechanism. With
 604 these cone-shaped edges, the soil volume lifted by the anchor is reduced, and
 605 consequently a slightly lower (better) upper-bound solution is obtained, as
 606 shown in Fig. 11.

607 One possible cause for the above mismatch is that no subdivision step is

608 performed for this anchor problem, as all velocity jumps are below the toler-
 609 ance for triggering h adaptivity (i.e., $\Delta v_{tol} \leq v_0$, where v_0 in this example is
 610 the anchor velocity). Accordingly, we lowered the tolerance Δv_{tol} to $0.1v_0$ to
 611 explore whether the Murray and Geddes (1987)'s solution can be recovered
 612 by refining the mesh. Figure 12 compares the optimized collapse mechanism
 613 based on the initial mesh and the refined one. It can be seen that because
 614 the initial mesh produces a uniform velocity field across elements (i.e., ve-
 615 locity jumps between elements are zero), only the blocks at the boundary
 616 of the failing soil volume are subdivided. In other words, this subdivision
 617 is based exclusively on the velocity jumps between moving blocks (i.e., the
 618 blocks $OBDA$ and $OCDB$ in Fig. 12(a)) and the assumed stationary region.
 619 Whereas a discontinuity passing the nodes D and O and insects the edge AB
 620 could lead to an improved mechanism, Figure 12(b) shows that the proposed
 621 subdivision strategy adds new discontinuities that intersect with the failure
 622 plane $ABCD$ and do not help in forming a more critical mechanism. As a

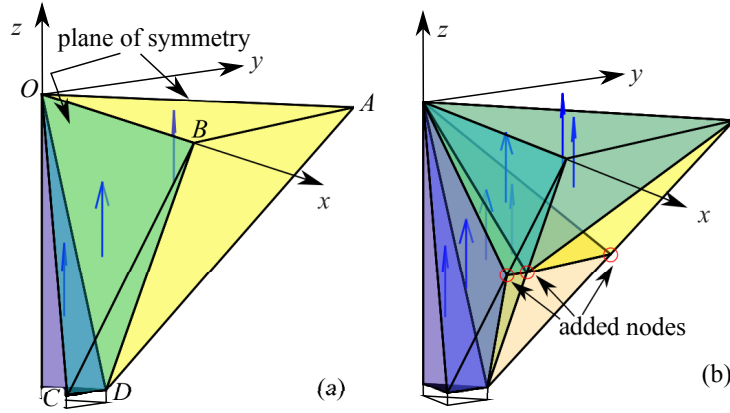


Figure 12: Collapse mechanism after optimizing nodal positions: (a) initial mesh; (b) after first subdivision.

consequence, the elements and discontinuities introduced in the refined mesh do not alter the optimized collapse mechanism or the computed limit load.

7.3. Rectangular wall in cohesionless soil

As a final example, the r - h adaptive approach is applied to compute the limit load on a rectangular retaining wall in cohesionless soil. The problem is illustrated in Fig. 13(a), which also depicts the collapse mechanism assumed by Soubra and Regenass (2000). The width and height of the wall are denoted by b and h , respectively, and the wall is assumed to move laterally into the soil (passive condition). The passive force on the wall at collapse can be expressed as

$$F = \frac{1}{2}\gamma b h^2 K_{\gamma p} \quad (18)$$

where $K_{\gamma p}$ is the so-called passive earth pressure coefficient. Generally, the passive resistance depends on the mode of wall movement, and in particular whether it translates, rotates, or moves with combined translation and

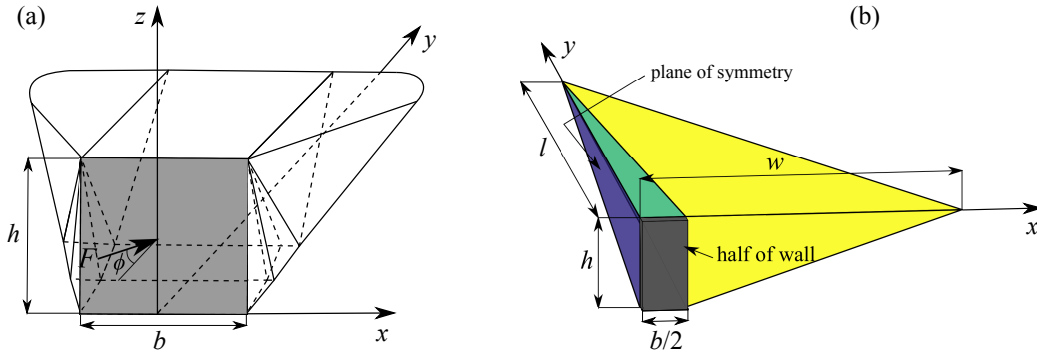


Figure 13: Collapse mechanism for passive failure of a rectangular retaining wall in cohesionless soil: (a) truncated multi-block mechanism (adapted from Soubra and Regenass (2000)); (b) initial mesh assumed in the r - h adaptive approach.

636 rotation (Widuliński et al., 2011). Here, only translational movement is con-
637 sidered.

638 This example models cohesionless soils in the same manner as those dis-
639 cussed earlier. The unit weight γ is assumed to be 1, and both the wall width
640 b and height h are assumed to be 2. Unlike the previous two examples, the
641 perfectly rough interface between soils and the retaining wall is modeled as
642 a velocity discontinuity, whose jump condition is characterized by friction
643 angle ϕ . This change is made in accordance with the assumption in the work
644 of Soubra and Regenass (2000), thus enabling a direct comparison.

645 Figure 13(b) shows the starting mesh used to initiate the computation.
646 Considering that $x = 0$ is a plane of symmetry, only half of the wall is
647 modeled. The geometric variables l and w are initially assumed to be $2h$,
648 and are adjusted to $l = 3.53h$ and $w = 7.20h$, for all values of ϕ , by the
649 sequential optimization discussed in Section 7.1. In this case the direction of
650 the force F is not horizontal but inclined at an angle ϕ with respect to the
651 direction normal to the wall (see Fig. 13(a)).

652 Figure 14 compares the $K_{\gamma p}$ values computed in this study and to those
653 assessed using the analytical solution proposed by Soubra and Regenass
654 (2000). The analytical solution corresponds to the failure mechanism shown
655 in Fig. 13(a), which consists of multiple blocks truncated by portions of cir-
656 cular cones. The methods provide very close results for small friction angles.
657 As the friction angle increases, the r - h adaptive approach gives lower (better)
658 estimates of the limit load.

659 The collapse mechanisms assessed through the r - h adaptive approach for
660 large and small friction angles are presented in Fig. 15. While both cases start

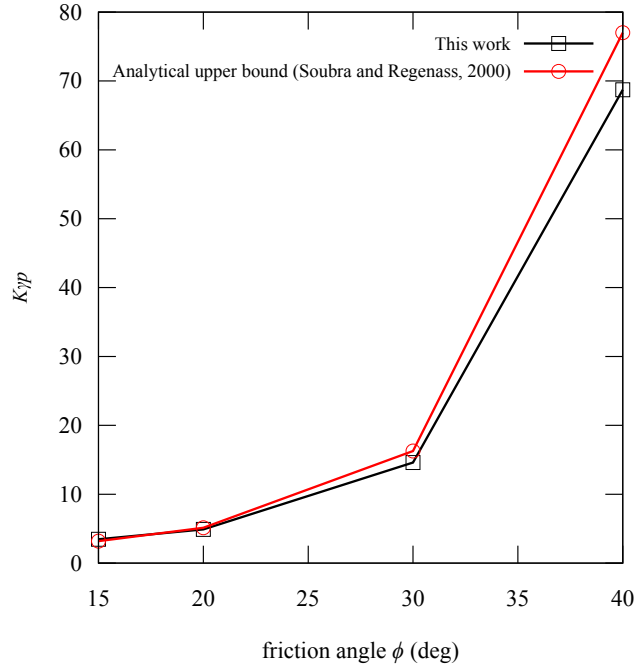


Figure 14: Comparison of the passive earth pressure coefficients $K_{\gamma p}$ computed by this work and the analytical solution of Soubra and Regenass (2000).

661 from the same initial mechanism, the proposed adaptive approach allows for
662 the mechanism to extend to greater depth and horizontal distance as the
663 friction angle grows, as in the solution of Soubra and Regenass (2000).

664 8. Discussion

665 Table 2 summarizes the computational cost of the r - h adaptive approach
666 for the three examples considered. The table includes the number of rigid
667 blocks (N_B), the number of nodal position components subjected to op-
668 timization (DOF), wall-clock times, and the CPU times required to run
669 MOSEK. Such information is organized for both the initial mesh configura-
670 tion and those after h adaptivity steps. For these examples, the r - h adaptive

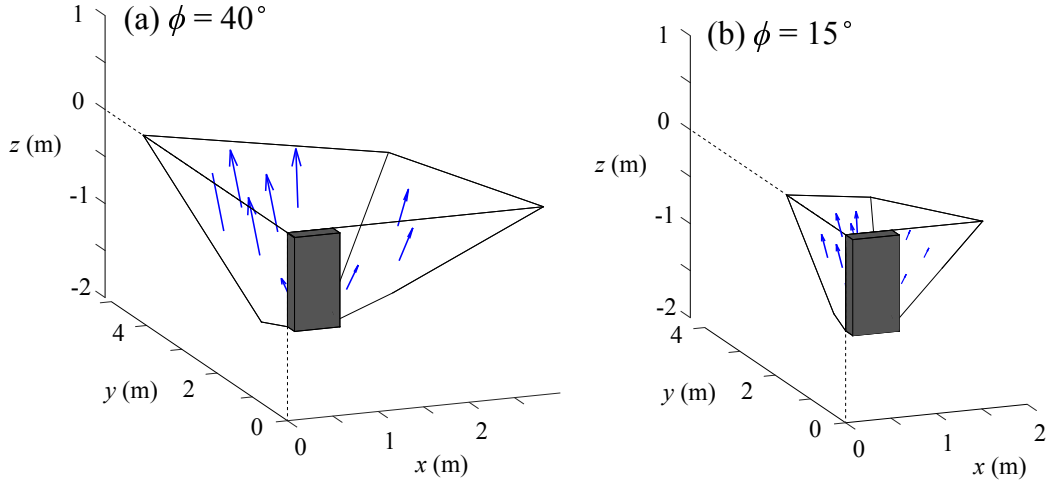


Figure 15: Collapse mechanisms computed with the r - h adaptive approach: (a) $\phi = 40^\circ$; (b) $\phi = 15^\circ$.

Table 2: Computational cost for the three examples.

Example	subdivision	N_B	N_{DOF}	wall-clock time (s)	MOSEK time (s)
square footing	0	3	3	0.8 to 1.2	0.1 to 0.3
	1	5 to 6	5 to 7	2.0 to 4.3	0.4 to 0.6
	2	11 to 14	11 to 18	4.4 to 21.2	0.9 to 4.2
retaining wall	0	3	2	1.7 to 2.1	0.6 to 0.7
	1	10 to 11	10 to 14	1.8 to 9.4	0.4 to 1.6
square anchor	0	3	2	1.8 to 2.2	0.7 to 0.8

671 approach displays promising computational efficiency. The maximum com-
672 putation time is no more than 30 seconds, observed in the square footing
673 case. The fact that the maximum MOSEK running time is around 4 seconds
674 suggests that the approach could also be accelerated by formulating a more
675 efficient strategy to solve the non-linear optimization problem of Eq. (12),

676 rather than using the FMINCON solver available in MATLAB.

677 While the above computation times are promising, they are not yet suf-
678 ficiently small to enable highly efficient sequential kinematic analysis for 3D
679 applications, one of the underlying objectives of this work. The computa-
680 tional demands of the proposed approach can be traced to the fact that a
681 forward numerical differentiation is employed to compute the gradient of the
682 objective function in the non-linear optimization. In other words, the ob-
683 jective function (the SOCP problem of Eq. 10) is called $DOF + 1$ times
684 to obtain the gradient. Therefore, computing the gradient consumes a sig-
685 nificant amount of time when DOF , corresponding to the number of nodal
686 positions, becomes large. To improve the computational efficiency, a future
687 refinement of the current work could be to approximate rather than directly
688 compute the gradient. For instance, the objective function can be linearized
689 with respect to its unknowns (cf. Hambleton and Sloan, 2013; Milani and
690 Lourenço, 2009), thus rendering an approximated but analytical form of the
691 gradient.

692 Compared with previous works on r adaptivity, an important contribu-
693 tion of this work is the adaptive subdivision, which automatically changes the
694 topological connectivity of blocks based on velocity jumps between blocks.
695 Table 3 summarizes the computed limit loads under initial mesh configura-
696 tion and after subsequent subdivisions. It can be seen that the calculated
697 upper bounds significantly decrease as the mesh is gradually refined, thus
698 suggesting that element subdivision based on the velocity jump is effective
699 in improving the collapse mechanism. Table 3 reports the computed results
700 only up to 2 subdivisions. The reason is that further mesh refinements only

Table 3: Limit loads computed by the initial mesh and after element subdivision (all reported values are obtained after optimizing nodal positions).

Numerical example	ϕ (°)	initial mesh	1st subdivision	2nd subdivision
square footing $N_{\gamma s}$ or N_{cs}	35	718.6	159	100.2
	30	150.5	50.9	35.0
	25	41.9	17.9	13.2
	20	13.6	6.6	5.4
	15	4.8	2.4	1.8
	0	8.3	6.7	6.4
retaining wall $K_{\gamma p}$	40	699.7	68.7	-
	30	23.9	14.6	-
	20	5.8	4.9	-
	15	3.5	3.5	-

lead to marginal improvement on the collapse loads. For example, additional mesh refinement only decreases the computed $N_{\gamma s}$ value in the square footing example by less than 5%, while the coefficient $K_{\gamma p}$ in the retaining wall problem remains unchanged even more subdivisions are performed.

The fact r - h adaptive approach eventually reaches a limit of no improvement can be attributed to two reasons. First, as demonstrated explicitly in Section 7.2, the proposed subdivision scheme does not always lead to an improvement in the solution. Indeed, the development of a more sophisticated subdivision strategy is an matter for future investigation. Such future algorithms can be devised by (1) identifying other useful indicators that flag the regions to be refined and (2) devising effective methods to subdivide elements so that discontinuities can be added at strategic locations. The second reason is that the algorithm used to solve the non-linear optimization

714 is only a local optimizer, and the solution is susceptible to being trapped at
715 a point that is a local rather than global optimum. The likelihood of this
716 occurring increases as elements are subdivided, since the the number of
717 unknowns (nodal positions) handled by the optimization is higher. As a part
718 of future work, global optimization techniques (e.g., genetic algorithm) can
719 be employed to resolve this potential limitation.

720 The initial mesh used as a starting guess in the proposed algorithm also
721 plays a significant role in the accuracy of the computed solution and whether
722 or not a global minimum can be attained. As made evident in the results
723 shown in Table 1, 3D DLO with a coarse grid may provide a reasonable
724 estimate of the limit loads at low cost, thus representing an encouraging
725 approach to systematically define initial meshes that can subsequently be
726 refined using the approach proposed in this work.

727 9. Conclusions

728 We propose an r - h adaptive kinematic approach for computing collapse
729 mechanisms and limit loads in 3D problems. Considering a velocity field
730 consisting of rigid elements (blocks) separated by zero-thickness velocity dis-
731 continuities, this method progressively improves the collapse mechanism and
732 bound on the limit load by successively adjusting the element nodal positions
733 (r adaptivity) as well as the element number and connectivity (h adaptivity).
734 Examination of the proposed technique through examples shows that when
735 the optimal mechanism is relatively simple, satisfactory limit loads can be
736 obtained solely by optimizing nodal positions (i.e., the locations of velocity
737 discontinuities), even if a simple mesh is assumed. However, when the op-

738 timal collapse mechanism becomes more intricate, adding discontinuities at
739 critical locations becomes crucial for the performance of r adaptivity. The
740 subdivision scheme proposed in this work automatically splits existing ele-
741 ments with velocity jumps greater than a specified threshold, adding new
742 elements so that velocity jumps can be further reduced through r adap-
743 tivity. This approach allows for the initiation of calculations from a very
744 simple mesh to which new discontinuities are progressively added at critical
745 locations, a paradigm that gives demonstrably high efficiency and may yield
746 higher efficiencies with future refinements.

747 To further speed up computations and enable efficient sequential kine-
748 matic analysis, wherein a full process of deformation is simulated through
749 a series of kinematic limit analysis computations, the proposed method can
750 be improved by pursuing alternatives to solving the non-linear optimiza-
751 tion of Eq. (12), devising more effective subdivision schemes, and developing
752 a systematic means of defining the initial mesh. These future refinements
753 represent important steps towards efficiently simulating large deformation
754 problems, especially those involving cohesionless soils, that are extremely
755 challenging to model by any other means.

756 10. Acknowledgements

757 Aspects of this work were supported by the Australian Research Council
758 through the award of Discovery Early Career Researcher Award (DE160100328)
759 and by the National Science Foundation (CMMI-1742849). This support is
760 gratefully acknowledged.

761 **References**

- 762 Antão, A. N., Santana, T. G., Vicente da Silva, M., da Costa Guerra, N. M.,
763 2011. Passive earth-pressure coefficients by upper-bound numerical limit
764 analysis. *Canadian Geotechnical Journal* 48 (5), 767–780.
- 765 Borges, L., Feijó, R., Zouain, N., 1999. A directional error estimator for
766 adaptive limit analysis. *Mechanics Research Communications* 26 (5), 555–
767 563.
- 768 Borges, L., Zouain, N., Costa, C., Feijó, R., 2001. An adaptive approach
769 to limit analysis. *International Journal of Solids and Structures* 38 (10),
770 1707–1720.
- 771 Chen, W.-F., 1975. *Limit Analysis and Soil Plasticity*. Elsevier.
- 772 Ciria, H., Peraire, J., Bonet, J., 2008. Mesh adaptive computation of upper
773 and lower bounds in limit analysis. *International journal for numerical*
774 *methods in engineering* 75 (8), 899–944.
- 775 Davis, E. H., 1968. Theories of plasticity and the failure of soil masses. In:
776 Lee, I. K. (Ed.), *Soil Mechanics: Selected Topics*. Butterworths, Sydney,
777 pp. 341–380.
- 778 Davis, E. H., Booker, J. R., 1971. The bearing capacity of strip footing
779 from the standpoint of plasticity theory. In: *Proc. the First Australia-New*
780 *Zealand Conference on Geomechanics*.
- 781 Drescher, A., 1991. *Analytical Methods in Bin-Load Analysis*. Vol. 36 of
782 *Developments in Civil Engineering*. Elsevier.

783 Drescher, A., Detournay, E., 1993. Limit load in translational failure mech-
784 anisms for associative and non-associative materials. *Géotechnique* 43 (3),
785 443–456.

786 Dunne, H. P., Martin, C. M., 2017. Capacity of rectangular mudmat foun-
787 dations on clay under combined loading. *Géotechnique* 67 (2), 168–180.

788 Godwin, R., O’Dogherty, M., 2007. Integrated soil tillage force prediction
789 models. *Journal of Terramechanics* 44 (1), 3–14.

790 Griffiths, D. V., Marquez, R. M., 2007. Three-dimensional slope stability
791 analysis by elasto-plastic finite elements. *Géotechnique* 57 (6), 537–546.

792 Hambleton, J. P., 2017. Earthmoving through the lens of geotechnical engi-
793 neering. In: 6th International Young Geotechnical Engineers Conference
794 (iYGEC6), Seoul, Korea, Sept. 17-22.

795 Hambleton, J. P., Drescher, A., 2012. Approximate model for blunt objects
796 indenting cohesive-frictional materials. *International Journal for Numerical
797 and Analytical Methods in Geomechanics* 36 (3), 249–271.

798 Hambleton, J. P., Sloan, S. W., 2013. A perturbation method for optimiza-
799 tion of rigid block mechanisms in the kinematic method of limit analysis.
800 *Computers and Geotechnics* 48, 260–271.

801 Hambleton, J. P., Sloan, S. W., 2016. A simplified kinematic method for 3D
802 limit analysis. *Applied Mechanics and Materials* 846, 342–347.

803 Hambleton, J. P., Stanier, S. A., White, D. J., Sloan, S. W., 2014. Modelling

804 ploughing and cutting processes in soils. *Australian Geomechanics* 49 (4),
805 147–156.

806 Hawksbee, S., Smith, C., Gilbert, M., 2013. Application of discontinuity
807 layout optimization to three-dimensional plasticity problems. *Proc R Soc*
808 *A* 469 (20130009).

809 He, L., Gilbert, M., 2016. Automatic rationalization of yield-line patterns
810 identified using discontinuity layout optimization. *International Journal of*
811 *Solids and Structures* 84, 27–39.

812 Hettiaratchi, D. R. P., Reece, A. R., 1974. The calculation of passive soil
813 resistance. *Géotechnique* 24, 289–310.

814 Johnson, D., 1995. Yield-line analysis by sequential linear programming. *In-*
815 *ternational Journal of Solids and Structures* 32 (10), 1395–1404.

816 Kashizadeh, E., Hambleton, J. P., Stanier, S. A., 2014. A numerical approach
817 for modelling the ploughing process in sands. In: *Proc. 14th International*
818 *Conference of the International Association for Computer Methods and*
819 *Advances in Geomechanics*, Kyoto, Japan, Sept. 22-25. pp. 159–164.

820 Kong, D., Martin, C. M., Byrne, B. W., 2017. Sequential limit analysis of
821 pipe–soil interaction during large-amplitude cyclic lateral displacements.
822 *Géotechnique* 68 (1), 1–12.

823 Krabbenhoft, K., Karim, M. R., Lyamin, A. V., Sloan, S. W., 2012. Asso-
824 ciated computational plasticity schemes for nonassociated frictional mate-
825 rials. *International Journal for Numerical Methods in Engineering* 90 (9),
826 1089–1117.

- 827 Krabbenhøft, K., Lyamin, A. V., Sloan, S. W., 2008. Three-dimensional
828 mohr–coulomb limit analysis using semidefinite programming. *Communi-*
829 *cations in Numerical Methods in Engineering* 24 (11), 1107–1119.
- 830 Lyamin, A. V., Salgado, R., Sloan, S. W., Prezzi, M., 2007. Two-and three-
831 dimensional bearing capacity of footings in sand. *Géotechnique* 57 (8),
832 647–662.
- 833 Lyamin, A. V., Sloan, S. W., 2002a. Lower bound limit analysis using non-
834 linear programming. *International Journal for Numerical Methods in En-*
835 *gineering* 55 (5), 573–611.
- 836 Lyamin, A. V., Sloan, S. W., 2002b. Upper bound limit analysis using lin-
837 ear finite elements and non-linear programming. *International Journal for*
838 *Numerical and Analytical Methods in Geomechanics* 26 (2), 181–216.
- 839 Lyamin, A. V., Sloan, S. W., Krabbenhøft, K., Hjiaj, M., 2005. Lower bound
840 limit analysis with adaptive remeshing. *International Journal for Numeri-*
841 *cal Methods in Engineering* 63 (14), 1961–1974.
- 842 Martin, C. M., 2005. Exact bearing capacity calculations using the method
843 of characteristics. In: *Proc. IACMAG. Turin.* pp. 441–450.
- 844 Martin, C. M., 2011. The use of adaptive finite-element limit analysis to
845 reveal slip-line fields. *Géotechnique Letters* 1 (4-6), 23–29.
- 846 Martin, C. M., Makrodimopoulos, A., 2008. Finite-element limit analysis of
847 mohr–coulomb materials in 3D using semidefinite programming. *Journal*
848 *of Engineering Mechanics* 134 (4), 339–347.

- 849 Mary, B. C. L., Maillot, B., Leroy, Y. M., 2013. Deterministic chaos in fric-
850 tional wedges revealed by convergence analysis. *International Journal for*
851 *Numerical and Analytical Methods in Geomechanics* 37 (17), 3036–3051.
- 852 Michalowski, R., Drescher, A., 2009. Three-dimensional stability of slopes
853 and excavations. *Géotechnique* 59 (10), 839–850.
- 854 Michalowski, R. L., 2001. Upper-bound load estimates on square and rect-
855 angular footings. *Géotechnique* 51 (9), 787–798.
- 856 Milani, G., 2015. Upper bound sequential linear programming mesh adapta-
857 tion scheme for collapse analysis of masonry vaults. *Advances in Engineer-*
858 *ing Software* 79, 91–110.
- 859 Milani, G., Lourenço, P. B., 2009. A discontinuous quasi-upper bound limit
860 analysis approach with sequential linear programming mesh adaptation.
861 *International Journal of Mechanical Sciences* 51 (1), 89–104.
- 862 Mosek, 2015. The mosek optimization toolbox for MATLAB manual. Version
863 7.1 (Revision 28), 17.
- 864 Muñoz, J. J., Bonet, J., Huerta, A., Peraire, J., 2009. Upper and lower
865 bounds in limit analysis: adaptive meshing strategies and discontinu-
866 ous loading. *International Journal for Numerical Methods in Engineering*
867 77 (4), 471–501.
- 868 Muñoz, J. J., Hambleton, J., Sloan, S. W., 2018. R-adaptivity in limit anal-
869 ysis. In: Barrera, O., Cocks, A., Ponter, A. (Eds.), *Advances in Direct*
870 *Methods for Materials and Structures*. Springer International Publishing,
871 Cham, pp. 73–84.

872 Murray, E. J., Geddes, J. D., 1987. Uplift of anchor plates in sand. *Journal*
873 *of Geotechnical Engineering* 113 (3), 202–215.

874 Nocedal, J., Wright, S. J., 2006. *Numerical optimization*. Springer.

875 Shield, R. T., Drucker, D. C., 1953. The application of limit analysis to
876 punch-indentation problems. *Journal of Applied Mechanics* 20, 453–460.

877 Sloan, S. W., 1995. Limit analysis in geotechnical engineering. In: Haberfield,
878 C. M. (Ed.), *Proc. Ian Boyd Donald Symposium on Modern Developments*
879 *in Geomechanics*, Melbourne, Australia, June 7.

880 Sloan, S. W., 2013. Geotechnical stability analysis. *Géotechnique* 63 (7), 531.

881 Smith, C., Gilbert, M., 2007. Application of discontinuity layout optimization
882 to plane plasticity problems. *Proc R Soc A* 463, 2461–2484.

883 Soubra, A.-H., Regenass, P., 2000. Three-dimensional passive earth pressures
884 by kinematical approach. *Journal of Geotechnical and Geoenvironmental*
885 *Engineering* 126 (11), 969–978.

886 Sturm, J. F., 2002. Implementation of interior point methods for mixed
887 semidefinite and second order cone optimization problems. *Optimization*
888 *Methods and Software* 17 (6), 1105–1154.

889 The MathWorks, Inc, 2018. *Optimization ToolboxTM User’s Guide*. The
890 MathWorks, Inc.

891 Vicente da Silva, M., Antão, A. N., 2008. Upper bound limit analysis with a
892 parallel mixed finite element formulation. *International Journal of Solids*
893 *and Structures* 45 (22), 5788–5804.

894 Widuliński, Ł., Teichman, J., Kozicki, J., Leśniewska, D., 2011. Discrete
895 simulations of shear zone patterning in sand in earth pressure problems
896 of a retaining wall. International Journal of Solids and Structures 48 (7),
897 1191 – 1209.

898 Wörden, F. T., Achmus, M., 2013. Numerical modeling of three-dimensional
899 active earth pressure acting on rigid walls. Computers and Geotechnics
900 51 (Supplement C), 83 – 90.