## Graphical Abstract

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## Highlights

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- Workflow for simulation of static and dynamic image analysis is presented.
- Particle characteristics for several solids types are determined.
- Influence image analysis methods on shape factors are shown.
- A correlation for 3D sphericity from 2D shape factors is derived.


# Evaluation of the Predictive Power of 2D Particle Imaging for 3D Characteristics in Bulk Material Analysis 

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#### Abstract

Particle size and shape characteristics are commonly measured with twodimensional (2D) imaging techniques, two of which are static or dynamic imaging techniques. These 2D particle characteristics need to be applied to particulate processes where they model three-dimensional (3D) processes. The correlation between 2D and 3D particle characteristics is therefore necessary, but the knowledge is still limited to either mathematically simple shapes or a certain set of investigated bulk solids.

A particle dataset consisting of six bulk solids measured with X-ray microscopy was used to simulate the results of 2D imaging techniques to create a database to test the correlation between sets of particle characteristics. The dataset thus created offers the possibility to study the correlation between characteristic values and robustly predict the 3D properties of bulk solids measured with 2D measurement techniques. It is found that the form factor, the square of circularity, is a good predictor of Wadell's sphericity, while the correlation can be improved by including additional 2D characteristics, namely convexity and the ratio of bounding circles.


Keywords: imaging techniques, static image analysis, dynamic image analysis, circularity, sphericity, shape factors, equivalent particle size, particle characteristics, correlation

## 1. Introduction

The characterization of particles regarding size and shape is essential for most particulate processes. Advances in measuring techniques have made the tomographic measurement of bulk solids and resulting particlediscrete datasets possible, enabling new methods of analyzing, e.g., separation processes [1,2].

However, tomographic measurement is a time-consuming and costly process, so the characterization of bulk solids in everyday industrial and laboratory applications is mostly done with other well-established techniques. For the measurement of particle size and shape in orders of magnitude from $1 \mu \mathrm{~m}$ to 10 mm , static and dynamic image analysis are widely used, and have often replaced traditional sieve analysis [3, 4, 5]. Furthermore, inline particle measurements are becoming more abundant in research and industry [6, 7].

Wadell introduced the concept of sphericity to account for a particle sedimentation velocity deviating from the sedimentation velocity of a sphere [8, 9]. It has since been used by many researchers and practitioners to represent particle shape as a single value. But Wadell recognized that the true sphericity for single particles might be hard to come by - it was even deemed unmeasurable by peers [10] - so he proposed the measurement of the projection of a particle at rest and an alternative definition for sphericity from it (Eq. 14).

The classical approach by Zingg to classify particles into shape categories by the ratios of their principal dimensions (elongation and flatness) is still widely in use and has been recently implemented in a particle shape analysis tool [11, 12]. 2D aspect ratios, along with circularity and convexity, are recognized in the literature as meaningful shape descriptors [13].

Since Wadell, many people have investigated how 2D imaging techniques may accurately describe the "true", 3D particle shape [14?, 15, 16, 17, 18]. In many ways, this study tries to retrace the steps of Bagheri et al. [19], who compared computed tomography measurements with projection images to find correlations to accurately describe 3D shape. Whereas before a particle's three principal dimensions (length, width, and thickness) were defined as perpendicular to each other, with length being the dimension between the two points on the particle furthest from each other, the authors propose the determination from the two projections with minimum (for thickness) and maximum areas (for width and thickness). Their
results are interesting, while lacking statistical robustness because of the small sample size.

Recent developments include the prediction of 3D particle shapes from 2D images by the use of neural networks [20]. This happens in recognition of the approach of capturing single particles from multiple angles to describe the 3D particle shape [21, 22]. The other approach is to quantify particle shape accurately only in the statistical sense by measuring enough particles to have a good estimate of the mean particle shape of a given bulk solid [13].

In this study, we take the second approach by asking how well 2D descriptors can describe 3D particle shape. We start with an expansive dataset of 3D particles provided by the PARROT particle database ${ }^{1}$ and simulate the results of both static and dynamic image analysis with the intent of finding suitable correlations for both methods.

## 2. Materials and Methods

### 2.1. Particle Datasets

### 2.1.1. Acquisition

The solids particle data used in this study was prepared previously for the stated purpose of providing reference 3D datasets. A methodology was developed to produce isolated, i.e., non-touching, particles in a wax matrix [23, 24]. Tomographic reconstruction of X-ray microscopy measurements of these wax matrices offers the possibility to easily segment and extract the single 3D particles. The particle data is available in the form of the original reconstructed tomography stacks as well as single particle surfaces in STL format in the dedicated particle database PARROT [25].

VTK files that represent cropped ROIs for every particle from the tomographic reconstructions were used to recalculate STL meshes for the particles, as some STL surfaces in the PARROT dataset were not watertight, which would have led to problems in later analysis. The STL data used in this study is available in the supplementary data.

Table 1 gives an overview of the six solids of which particle surface data has been used. They are typically in a particle size range between $50 \mu \mathrm{~m}$ to

[^0]$300 \mu \mathrm{~m}$. The X-ray microscopy measurements were performed for a final voxel size, i.e., edge length, of $2 \mu \mathrm{~m}$.

Table 1: Used Particle Systems, provided in the PARROT particle database [25]

| type | production process | particle size | particles |
| :--- | :--- | :--- | :---: |
| aluminium <br> oxide <br> dolomite | crushing | $55 \mu \mathrm{~m}$ to $200 \mu \mathrm{~m}$ | 1571 |
| calcination and |  |  |  |
| crushing <br> glass | $90 \mu \mathrm{~m}$ to $200 \mu \mathrm{~m}$ | 642 |  |
| spray drying | $150 \mu \mathrm{~m}$ to $300 \mu \mathrm{~m}$ | 602 |  |
| mica | dry milling <br> comminution and | $55 \mu \mathrm{~m}$ to $200 \mu \mathrm{~m}$ | 1271 |
|  | magnetic separa- | 400 m to 300 | 415 |
| quartz | tion <br> crushing | $<200 \mu \mathrm{~m}$ | 1656 |

### 2.1.2. Description

The properties of the six solids (cf. Table1) are shown in Fig.1. From the plot of sphericity $\psi_{W a}$ over volume-equivalent diameter (Eqs. 8 and 2) in Fig.1a, it can be seen that four solids-quartz, limestone, dolomite, and aluminum oxide-are clustered in the same area with relatively high sphericity values of $\psi_{\mathrm{Wa}}>0.5$. The soda-lime glass particles are the largest and also have the highest sphericity values. The high sphericity values can be traced to the production process by spray drying, resulting in mostly spherical shapes. In contrast, the mica particles show very low sphericities.

The maximum sphericity value of $\psi_{\mathrm{Wa}} \approx 0.92$ stems from the conversion of the particle volumes from voxel representation to a triangular surface. The marching cubes algorithm interpolates between the edges of the voxels to smooth the surface, depending on the number and configuration of adjacent solid voxels [26,27]. The resulting error will be $8 \%$ to $9 \%$ [28]. In comparison, the error in sphericity determination from the voxel representation for a sphere would be $>30 \%$, because of the greatly exaggerated surface area.

In Fig.1d, the particles are plotted along two aspect ratios, flatness $t / w$ and elongation $w / l$, which makes classification according to particle shape


Figure 1: Properties of the six particle datasets
possible. $l, w$, and $t$ are the three main dimensions of a particle: length, width, and thickness, respectively, defined by the aligned bounding box (cf. section2.3.2). The plot was first introduced by Zingg and later developed by Krumbein and Janoo [11, 29, 30]. Fig.1c serves as an explanation, also showing isolines for sphericity, though Krumbein's sphericity definition is used, cf. Eq.11. Alternative descriptors for the particle shape groups "disc," "cubic," and "rod" are "oblate," "compact," and "prolate," respectively [31].

Soda-lime glass spheres are expectedly clustered at values close to one for both aspect ratios, while the majority of particles of the other solids are mostly compact and could be classified as cubic and slightly rod- or disc-shaped, depending on their particular flatness or elongation values.

In contrast, the mica particles are very flat and may be classified as discand blade-like.

Fig. 2 provides an example for each of the four categories according to the Zingg classification chart. The examples also serve to give an impression of what the different solids look like. While most of the limestone and quartz particles can be classified as compact/cubic, the two particles shown in Figs.2a and 2d can be clearly identified as belonging to their respective categories of disc- and rod-like. The soda-lime glass particles are mostly near-perfect spheres, resulting in the aforementioned high sphericity values. The mica is mostly flaky in nature, resulting in very low flatness values, an effect that can be predicted from the plot of width vs. thickness in Fig. 1b.


Figure 2: Examples from the datasets for particles belonging to the four shape categories of Figs. 1c and 1d with solid type in brackets

Because their properties are very similar, the group of quartz, limestone, dolomite, and aluminum oxide will be grouped as "compact particles" in section 3, while it will be instructive to see certain deviations for the
mica and soda-lime glass particles occur in the calculation of form factors because of their unique shape properties.

### 2.2. 2D Imaging Simulation

2.2.1. Static Image Analysis

Static image analysis, as defined by ISO 13322-1, involves image acquisition to determine particle size where the particles are not moving against the axis of the optical equipment [32]. If a particle is large enough that adhesion forces with respect to the surface it is resting on are negligible, the particle will orient itself in a position in which at least its longest dimension is measurable. Two possibilities for the simulation of static image analysis were calculated:

- alignment of the principal inertia vectors on the Cartesian axes and
- alignment of the particle in one of its stable resting positions.

3D manipulation of the provided STL files was done with the Python library trimesh, which, as the name implies, focuses on triangular meshes [33]. The trimesh package offers options for both the procedures named, in particular a method that returns a list of the most likely stable positions of a given mesh, containing both the necessary transform and the respective probability of the particle settling in this position. Any resting positions with a probability $p>0.1$ were used for further 2D analysis. Because highly spherical particles can easily have no positions of especially high resting probability, for each particle at least the two most probable positions were calculated. Fig. 3 gives an example of the stable positions of a particle and the resulting projections, in this case in $z$-direction, i.e., onto the $x y$ plane.

The imaging simulation involves getting the projection perpendicular to the plane that acts as the resting surface when calculating the stable position transforms $(x y)$. For the mesh aligned along its principal inertia vectors, the projection is calculated perpendicular to the plane that contains the two major inertia vectors: when considering the aligned particle in Fig. 6b, the projection would be in direction of the $x$ vector, onto the $y z$ plane.

The subsequent procedure involves a custom function that calculates the orthogonal projection of the triangular mesh onto a plane defined by a given normal. With a given plane normal, the particle is first rotated to
the correct position, and a projection transform is performed onto the $y z$ ( $x$-axis) plane (Fig. 4a). The projected triangles are then transformed into a single 2D polygon using the Python package Shapely [34]. Thus, a single contour is returned which can be used for further analysis. The relevant code can be found in the supplementary materials, see section 4.

In principle, the effects of image resolution may be investigated by scaling the projection and calculating a masked array that represents the pixel image. However, pixelization in this sense has only been used for the calculation of the enclosing and inscribed circles, cf. section 2.3.1.


Figure 3: Stable position of the particle shown in fig. 6 with the respective occurrence probabilities (figs. 3a, 3b, and 3c) and resulting projection silhouettes along $z$-axis (figs. 3d, $3 e$, and $3 f$ )

### 2.2.2. Dynamic Image Analysis

In contrast to static analysis, dynamic image analysis is concerned with the image acquisition and analysis of moving particles [35]. Particles are therefore imaged in random orientations, unless the flow is highly turbulent.

The procedure to produce a projection image is mostly the same as before, except that the particle is first rotated randomly. For every particle, three random orientations were used to produce projections, thereby increasing the number of simulated data points.

### 2.3. Particle Characteristics

The term particle characteristic as used in this text includes all parameters that can describe the size and shape of a particle. It comprises three subgroups: geometric properties, equivalent diameters and shape factors. Geometric properties can be directly measured from the 2D or 3D representation of a given particle. Equivalent diameters are typically diameters of the circle (2D) or sphere (3D) that share one of the geometrical properties of the particle. Finally, shape factors are mostly ratios of two different geometrical properties, one of which may be calculated from the particle's convex hull.

### 2.3.1. 2D Measures

These geometric properties can all be derived directly from the projection or section of a particle in any direction (Fig. 4b); therefore, they are applicable to all 2D imaging techniques, like static and dynamic image analysis.

In the current study, only the vector representation of the silhouette image is used. The accuracy of the calculated parameters therefore only depends on the resolution of the original 3D surface mesh and the marching cubes procedure with which it was produced from the voxel representations that themselves originated in the reconstructed tomography image stack. The contour is voxelized solely to simplify the calculation of bounding circles, enabling the use of standard Python libraries. Both pixelization and orthogonal projection images, as shown in Fig. 4a, offer possibilities for testing the effects of image resolution and roughness measurement, respectively.

Area and perimeter. Both the projection area $A_{\mathrm{p}}$ and the perimeter $P_{\mathrm{p}}$ are calculated by methods provided by the Shapely package, directly from the projection contour, as shown in Fig. 4b. Because of the inherent fractal behavior of many real solids' surfaces, the perimeter is much less robust than the projection area for smaller particles. Still, the effect of measurement resolution will be more pronounced in the determination of the (3D) surface area, where surface roughness comes more into effect than in the 2D case [36].


Figure 4: Illustration of methods for generation of 2D descriptors for particle shape

Convex Hull. The convex hull is determined using a method of the Shapely polygon object that contains the contour. For the convex hull, both area $A_{\mathrm{c}}$ and perimeter $P_{\mathrm{c}}$ are determined.

Feret Diameters. Minimum and maximum Feret diameters are determined by brute force: the projection contour is rotated in 500 steps between $0^{\circ}$ and $180^{\circ}$, and the boundaries in both axis directions are determined. The smallest measured distance between boundaries will be the minimum Feret diameter $x_{\mathrm{Fe} \text {, min }}$, while the largest distance will be the maximum Feret diameter $x_{\mathrm{Fe}, \max }$. The two measures, $x_{\mathrm{Fe}, \min }$ and $x_{\mathrm{Fe}, \max }$, are shown in Fig. 5a. As can be seen, the two Feret diameters are not necessarily at a right angle, which is why two additional Feret diameters are determined: $x_{\mathrm{Fe}, \min 90}$ and $x_{\mathrm{Fe}, \max 90}$, which are perpendicular to the $x_{\mathrm{Fe}, \max }$ and $x_{\mathrm{Fe}, \min }$, respectively.

The use of perpendicular Feret diameters serves two purposes. Firstly,


Figure 5: Illustration of of different definitions of Feret diameters
for static image analysis, the maximum and minimum Feret diameters will be very close to the length and width of a particle, respectively (Fig. 5b). Secondly, the (true) minimum Feret diameter $x_{\mathrm{Fe}, \mathrm{min}}$ and its perpendicular Feret diameter $x_{\mathrm{Fe}, \max 90}$ will, in most cases, be very close to the actual dimensions of the oriented bounding box, i.e., the bounding box of least area.

Minimum Enclosing Circle. The diameter of the minimum enclosing circle $d_{\text {ec }}$ belongs to the circle that has the least area while still containing the entire projection contour (Fig. 4f). While dedicated Python packages for the task of determining this measure exist, such as miniball, here, the computer vision library OpenCV was used [37].

For the calculation of $d_{\mathrm{ec}}$, the contour needs to be transformed into a array first, equivalent to a pixel representation (Fig. 4d). The pixelization is achieved with scikit-image, which contains the polygon method that generates pixel coordinates inside a given polygon.

To increase the accuracy of $d_{\mathrm{ec}}$ (and $d_{\mathrm{ic}}$ ), the contour coordinates are scaled up by a factor of 2 before pixelization, significantly affecting on the results of both the center coordinate of the circle as well as its radius. Further scale-up is not considered necessary, or even useful, because the original 3D mesh does not offer more resolution anyway.

Maximum Inscribed Circle. The determination of the diameter of the maximum inscribed circle $d_{\text {ic }}$ also requires a pixel representation of the contour. The method uses the Euclidean distance transform as implemented in scipy [38]. The transform calculates the distance of each object pixel from
the background (Fig. 4e). The pixel that contains the highest value after the transform will be the center of the maximum inscribed circle, while the corresponding pixel value will be $d_{\text {ic }} / 2$, i.e., the radius of the circle. The Euclidean distance transform is computationally inexpensive and is a relatively simple method for determining the maximum inscribed circle, as it transforms the problem from vector space to pixel space. This reduces the complexity of the problem significantly, albeit at the cost of being only as accurate as the pixel dimensions allow.

### 2.3.2. 3D Measures

Volume and Surface Area. Both volume and surface area are properties of the trimesh object that contains the particle mesh, so it is defined by functions already implemented by the package.

Specific Surface Area. A combination of volume and area, specific surface area is an important measure for all sorts of processes involving heat, moment, or mass transfer. It is defined as:

$$
\begin{equation*}
S_{\mathrm{V}}=\frac{S}{V} \tag{1}
\end{equation*}
$$

In contrast to most other particle properties, specific surface area will decrease with increasing volume.

Convex Hull. The convex hull is another property of the trimesh object, from which both volume $V_{\mathrm{c}}$ and surface area $S_{\mathrm{c}}$ can be calculated.

Aligned Bounding Box. A bounding box in this study defines the main dimensions of the particle. In this study, the aligned bounding box defines the length $l$, width $w$, and thickness $t$ to be the longest, intermediate, and shortest edge lengths. This approach is congruent with the definition of particle dimensions by Krumbein, who measured orthogonal lengths starting with the longest one found on the particle [29].

The aligned boudning box is created by tranformation of the particle so that its principal axes of inertia align with the cartesian dimension vectors (Fig. 6b). The necessary transform is again a property of the trimesh object containing the particle mesh. After the transformation, the bounding box, again, is a property of the trimesh object (Fig. 6d).

The definition of 3D particle dimensions in this way also makes it possible to directly compare measurements with static image analysis simulation results. When the maximum Feret $x_{\mathrm{Fe}, \max }$ and the perpendicular


Figure 6: Illustration of the two different definitions for bounding boxes, volumes given in axis units

Feret diameter $x_{\mathrm{Fe}, \min 90}$ (Fig. 5b) is used, they will be identical with length $l$ and width $w$ for the aligned particle (section 3.1). For stable positions, section $3.2, x_{\mathrm{Fe}, \max }$ should still reflect actual particle length $l$, while $x_{\mathrm{Fe}, \min 90}$ should differ somewhat.

Bagheri et al. favoured the use of uncorrelated Feret extrema for the determination of particle dimensions to reduce operator error [19]. However, with most modern measurement setups particle dimensions are seldom determined manually, and determination of a minimum Feret diameter for compact projections may still be difficult if done manually anyway.

Oriented Bounding Box. The oriented bounding box is again calculated by trimesh for a given particle mesh and represents the bounding box of least volume that still contains the whole mesh surface (Fig. 6c). The dimensions of the oriented bounding box are determined from the Cartesian coordinates after applying the inverse transform on the bounding box, since the oriented bounding box is likely to be at random angles toward the Cartesian axes, even if the particle was first aligned to its principal axes of intertia.

Fig. 7 shows comparisons of the dimensions of aligned bounding boxes and oriented bounding boxes for all investigated particles. The oriented bounding box has on average smaller dimensions than the aligned bounding box. The effect increases for the longer dimensions: length will mostly be smaller for the oriented bounding box, wheras there is a more random scatter for thickness.


Figure 7: Comparison of dimensions determined by aligned and oriented bounding boxes

On average, the oriented bounding box will be $14 \%$ smaller than the aligned bounding box for compact particles. In contrast, the oriented bounding box will only be $12 \%$ smaller for mica particles which, because
of their flat nature, should, in their aligned position, already be closer to the smallest box possible. Finally, soda-lime glass spheres have on average oriented bounding boxes that are only $5.5 \%$ smaller.

The aligned bounding box is preferred here over the oriented bounding box because of its congruence with Krumbein's definition and because the resulting dimensions could be found more easily by hand.

Bounding Spheres. The minimum bounding sphere again is a property of the mesh object defined by the trimesh library, so the diameter of the minimum enclosing sphere $d_{\text {es }}$ is determined in a single line of code. A visualization of both bounding spheres is found in Fig. 8.


Figure 8: Illustration of both minimum enclosing sphere and maximum inscribed sphere
The maximum inscribed sphere is approximated as the maximum inscribed circle in the 2D case. In both cases, the function distance_transform_edt from the scripy library [38] is used to calculate the Euclidean distance transform to find the pixel/voxel that is furthest from the particle surface. This maximum value will be the diameter of the maximum inscribed sphere $d_{\text {is }}$.

In order to perform the Euclidean distance transform, the surface mesh needs to be discretized into a voxel representation (Fig. 8c). The voxelization is also done with methods provided by trimesh, and, as with the 2D case, at a scale factor of 2 , which increases the accuracy of the diameter estimation significantly. Care must be taken to produce a filled voxel representation: most voxelization algorithms will only return solid voxels where the surface of the mesh touches. An extra step is involved to fill the hollow discretized surface with scipy's method binary_fill_holes.

### 2.3.3. Equivalent Diameters

Several properties in 2D and 3D can be compared to that of the idealized shapes, a circle in two and a sphere in three dimensions. In 3D, the diameter of a sphere can be calculated that has the same volume as that of the particle. This diameter will be called the volume-equivalent diameter:

$$
\begin{equation*}
x_{\mathrm{V}}=\sqrt[3]{\frac{6 V_{\mathrm{p}}}{\pi}} \tag{2}
\end{equation*}
$$

In the same sense, the diameter of the sphere that has the same surface area as that of the particle (surface-equivalent diameter) is:

$$
\begin{equation*}
x_{\mathrm{S}}=\sqrt{\frac{S_{\mathrm{p}}}{\pi}} \tag{3}
\end{equation*}
$$

In two dimensions, the particle properties volume $V_{\mathrm{p}}$ and surface area $A_{\mathrm{p}}$ reduce to projection properties, projection area $A_{\mathrm{p}}$ and perimeter $P_{\mathrm{p}}$. The diameter of the circle that has the same area as the projection area, the area-equivalent diameter, is:

$$
\begin{equation*}
x_{\mathrm{A}}=\sqrt{\frac{4 A_{\mathrm{p}}}{\pi}} \tag{4}
\end{equation*}
$$

Lastly, the perimeter-equivalent diameter is the diameter of the circle that has the same perimeter as that of the particle projection, defined as:

$$
\begin{equation*}
x_{\mathrm{P}}=\frac{P_{\mathrm{p}}}{\pi} \tag{5}
\end{equation*}
$$

### 2.3.4. Shape Factors

Shape factors are derived from two or three of the particle properties or equivalent diameters introduced above. All shape factors described below are dimensionless, which means they can be used to good effect to find correlations between 2D projections and 3D particle properties.

Length Ratios. Flatness $t / w$ and elongation $w / l$ have been used before in Fig. 1d to classify particles into shape categories.

In 2D, two more length ratios are used in this study. First the aspect ratio is defined as the ratio of minimum and maximum Feret diameter:

$$
\begin{equation*}
\mathrm{AR}=\frac{x_{\mathrm{Fe}, \min }}{x_{\mathrm{Fe}, \max }} \tag{6}
\end{equation*}
$$

As discussed before, the two Feret diameters often at an angle $\neq 90^{\circ}$. Because the 3D particle dimensions are defined by their bounding boxes, they are necessarily at a right angle to each other. It therefore makes sense to define an additional aspect ratio of perpendicular Feret diameters:

$$
\begin{equation*}
\mathrm{AR}_{90}=\frac{x_{\mathrm{Fe}, \min 90}}{x_{\mathrm{Fe}, \max }} \tag{7}
\end{equation*}
$$

Sphericity. Several shericity definitions exist, some of them fundamentally different from each other, but for all of them, the sphericity $\psi<1$ for particles deviating from a sphere.

The original definition of sphericity comes from Wadell for application on sedimentary particles [8]. Wadell defined sphericity as the ratio of the surface area of a sphere of equal volume as that of the particle to the actual surface area of the particle:

$$
\begin{equation*}
\psi_{\mathrm{Wa}}=\frac{S_{\mathrm{sp}}}{S_{\mathrm{p}}}=\left(\frac{x_{\mathrm{V}}}{x_{\mathrm{A}}}\right)^{2} \tag{8}
\end{equation*}
$$

$S_{\text {sp }}$ is the surface area of the sphere having the same volume as the particle.

Another sphericity definition is the ratio of the two bounding spheres, i.e., maximum inscribed sphere to minimum enclosing sphere [39]:

$$
\begin{equation*}
\psi_{\mathrm{bs}}=\frac{d_{\mathrm{is}}}{d_{\mathrm{es}}} \tag{9}
\end{equation*}
$$

Hofmann applies the concept of statistical entropy to the particle shape description [40]:

$$
\begin{equation*}
\psi_{\mathrm{Ho}}=\frac{1}{\ln (1 / 3)} \sum_{i=1}^{3} p_{i} \ln p_{i} \tag{10}
\end{equation*}
$$

where $p_{i}=\frac{d_{i}}{d_{1}+d_{2}+d_{3}}, d_{1}=l, d_{2}=w$, and $d_{3}=t$.
Hofmann's sphericity is supposed to be the most representative measure for the prediction of particle settling velocity [41].

Lastly, Krumbein defined a sphericity by comparing a given particle to a triaxial ellipsoid [29]. After determining the longest dimension of the particle, the second longest dimension perpendicular to the first is determined, with the third dimension being perpendicular to the other two.

In this sense, the three dimensions are equivalent to length $l$, width $w$, and thickness $t$ of the bounding box of the principally aligned particle, as described in section 2.3.2.

$$
\begin{equation*}
\psi_{\mathrm{Kr}}=\sqrt[3]{\frac{w t}{l^{2}}} \tag{11}
\end{equation*}
$$

Another definition for sphericity has been defined by Sneed and Folk as $\psi_{\mathrm{SF}}=\sqrt[3]{t^{2} /(w l)}$ [42], but will not be used in this study.

Circularity. "Circularity" is the name chosen according to the definitions of Wadell [9] for the 2D equivalent of sphericity, basically a "projection sphericity", sometimes also called "roundness" [43]. Like sphericity, circularity approaches a value of one for particles that closely resemble circular shapes and will decrease in value for particles becoming less compact.

The original circularity definition as ratio of perimeter of the areaequivalent circle to the actual projection perimeter is due to Wadell [9]. Wadell stressed that circularity and sphericity are fundamentally different from roundness in the sense that roundness is a mesoscopic measure and circularity is a macroscopic measure. In other words, circularity and sphericity show shape deviations, whereas roundness shows surface deviations.

$$
\begin{equation*}
\psi_{\mathrm{c}}=\frac{P_{\mathrm{c}}}{P_{\mathrm{p}}}=\frac{x_{\mathrm{A}}}{x_{\mathrm{P}}}=\sqrt{\frac{4 \pi A_{\mathrm{p}}}{P_{\mathrm{p}}^{2}}} \tag{12}
\end{equation*}
$$

The square of circularity $\psi_{c}$ is called the form factor and is equivalent to the "roundness" factor defined by Cox [44, 45, 46].

$$
\begin{equation*}
\mathrm{FF}=\frac{4 \pi A_{\mathrm{p}}}{P_{\mathrm{p}}^{2}} \tag{13}
\end{equation*}
$$

Because one early criticism of $\psi_{\mathrm{Wa}}$ was the difficulty of measurement, Wadell proposed more easily attainable circularity measure:

$$
\begin{equation*}
\psi_{\mathrm{c}, \mathrm{Wa}}=\frac{x_{\mathrm{A}, \text { stable }}}{d_{\mathrm{ec}}} \tag{14}
\end{equation*}
$$

In the above equation, $x_{\mathrm{A}, \text { stable }}$ is the diameter of a circle of equal projection area as that of a given particle at rest, i.e., lying on a surface in a stable
position. $d_{\mathrm{ec}}$ is, as per previous definition, the diameter of the minimum enclosing circle.

Another method of defining circularity is through both bounding circles, i.e., the radius of the maximum inscribed circle $d_{\mathrm{ic}}$ and the radius of the minimum enclosing circle $d_{\mathrm{ec}}$ :

$$
\begin{equation*}
\psi_{\mathrm{c}, \mathrm{bc}}=\frac{d_{\mathrm{ic}}}{d_{\mathrm{ec}}} \tag{15}
\end{equation*}
$$

Equation 15 is the square of the circularity definition by Riley [43].
Solidity. As a measure of concavity, a solidity factor $S_{x}$ can be calculated in both 2D and 3D. It compares the actual particle volume or projection area to its convex hulls. If there are no concavities, the solidity will be 1 and the particle or projection will be its own convex hull.

$$
\begin{align*}
& S_{\mathrm{x}, 3 \mathrm{D}}=\frac{V_{\mathrm{p}}}{V_{\mathrm{c}}}  \tag{16}\\
& S_{\mathrm{x}, 2 \mathrm{D}}=\frac{A_{\mathrm{p}}}{A_{\mathrm{c}}} \tag{17}
\end{align*}
$$

Convexity. Another measure for deviation from a convex object is the convexity, for which the symbol $C_{x}$ is used. It compares the surface of particle or projection directly to the convex hull.

$$
\begin{align*}
C_{x, 3 \mathrm{D}} & =\frac{S_{\mathrm{c}}}{S_{\mathrm{p}}}  \tag{18}\\
C_{\mathrm{x}, 2 \mathrm{D}} & =\frac{P_{\mathrm{c}}}{P_{\mathrm{p}}} \tag{19}
\end{align*}
$$

## 3. Results and Discussion

### 3.1. Aligned Projection

The aligned projection dataset is in many ways the simplest one and is used for verification of the analysis methods then used for the datasets of stable and dynamic projections. Because there is exactly one aligned projection for every particle, there are as many projections as particles in
the complete dataset of all solids, 6157. Because of the amount of particles, any effects observed are considered statistically relevant.

The total number of particle characteristics used for correlation is 49,25 comprising 3D, 24 comprising 2D measures and descriptors. Table 2 lists all particle characteristics, which have been grouped into certain categories like volume-related, circularity, etc.

These characteristics can now be used to calculate a correlation matrix as shown in Fig. 9a. Simply put, the Pearson correlation coefficient of each parameter is evaluated against every other parameter, resulting in a $49 \times 49$ grid containing the values of the coefficients. From the numbers on the grid, the specific characteristic can be determined with Table 2.

Values greater than zero will signify a positive (linear) correlation, whereas, if rarely, negative values will signify negative (linear) correlations. Extremely high correlations result from some expected pairs, like the equivalent diameters and their respective measure, or circularity (44) and form factor (45) - one is the square of the other.

For the geometric measures and equivalent diameters a clear dependency is visible by four dark blue rectangles that are formed. The brighter regions of less correlation are all in places of shape factors. A obvious exception from the rule is specific surface area (9), that decreases with increasing volume and therefore results in a band of negative correlation throughout the correlation matrix (Fig. 9b).

One correlation that is not necessarily expected is between Wadell's alternative circularity definition $\psi_{c, W a}$ (46, Eq. 14) and the bounding circles circularity $\psi_{c, b c}(47$, Eq. 15). As predicted by the correlation plot, there is a near perfect linear relationship, but between $\psi_{c, W a}^{2}$ and $\psi_{c, b c}$, as shown in Fig. 10.

Because the main focus of this study is the comparison of 2D with 3D particle characteristics, most of the correlation matrix is not strictly relevant. For this reason, only the upper right quadrant is shown for the other correlation matrices, as has been done for the larger matrix in Fig. 9c for the set of compact particles: quartz, limestone, dolomite, and aluminium oxide. Marked in red are characteristics pairs of very high correlation. Thresholds for a "high" correlation are set subjectively, as shape factors overall show much less correlation than geometric measures and their derived equivalent diameters. In Fig. 9c, some expected characteristics show high correlation like convex surface area (6) and (convex) projection area $(1,2)$, or their equivalent diameters: $x_{\mathrm{S}}(7)$ and $x_{\mathrm{S}, \mathrm{c}}(8)$ with $x_{\mathrm{A}}(3)$

Table 2: Particle characteristics used in the correlation matrices, Figs. 9, ...

| no. category name |  |  | symbol | dimensions | equation |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 3D characteristics |  |  |  |  |  |
| 1 | I | particle volume convex volume volume-equivalent diameter convex volume-eq. diameter | $V_{\mathrm{p}}$ | $L^{3}$ | - |
| 2 |  |  | $V_{\text {c }}$ | $L^{3}$ | - |
| 3 |  |  | $x_{\mathrm{V}}$ | L ${ }^{1}$ | 2 |
| 4 |  |  | $x_{\mathrm{V}, \mathrm{c}}$ | $\mathrm{L}^{1}$ | - |
| 5 | II | particle surface area convex surface area surface-equivalent diameter convex surface-eq. diameter volume-specific surface area | $S_{p}$ | L ${ }^{2}$ | - |
| 6 |  |  | $S_{\text {c }}$ | $\mathrm{L}^{2}$ | - |
| 7 |  |  | $x_{\text {S }}$ | $L^{1}$ | 3 |
| 8 |  |  | $x_{\text {S, } \mathrm{c}}$ | $\mathrm{L}^{1}$ | - |
| 9 |  |  | $S_{\text {V }}$ | $\mathrm{L}^{-1}$ | - |
| 10 | III | aligned length oriented length | $l$ | L ${ }^{1}$ | - |
| 11 |  |  | $l_{\text {oriented }}$ | $\mathrm{L}^{1}$ | - |
| 12 | IV | aligned width oriented width | $w$ | $\mathrm{L}^{1}$ | - |
| 13 |  |  | $w_{\text {oriented }}$ | $\mathrm{L}^{1}$ | - |
| 14 | V | aligned thickness oriented thickness | $t$ | L ${ }^{1}$ | - |
| 15 |  |  | $t_{\text {oriented }}$ | $\mathrm{L}^{1}$ | - |
| 16 | VI | min. enclosing sphere diameter max. inscribed sphere diameter | $d_{\text {es }}$ | $L^{1}$ | - |
| 17 |  |  | $d_{\text {is }}$ | $\mathrm{L}^{1}$ | - |
| 18 | VII | elongation flatness | $w / l$ | - | - |
| 19 |  |  | $t / w$ | - | - |
| 20 | VIII | Wadell's sphericity Krumbein's sphericity bounding spheres sphericity Hofmann's sphericity | $\psi_{\text {Wa }}$ | - | 8 |
| 21 |  |  | $\psi_{\mathrm{Kr}}$ | - | 11 |
| 22 |  |  | $\psi$ Wa | - | 9 |
| 23 |  |  | $\psi_{\text {Ho }}$ | - | 10 |
| 24 | IX | 3D solidity 3D convexity | $S_{x, 3 \mathrm{D}}$ | - | 16 |
| 25 |  |  | $C_{x, 3 D}$ | - | 18 |
| 2D characteristics |  |  |  |  |  |
| 26 | I | projection area convex projection area area-equivalent diameter convex area-eq. diameter | $A_{\mathrm{p}}$ | L ${ }^{2}$ | - |
| 27 |  |  | $A_{\text {c }}$ | $\mathrm{L}^{2}$ | - |
| 28 |  |  | $x_{\text {A }}$ | $\mathrm{L}^{1}$ | 4 |
| 29 |  |  | $x_{\text {A, }}$ | $L^{1}$ | - |
| 30 | II | projection perimeter convex projection perimeter perimeter-equivalent diameter convex perimeter-eq. diameter | $P_{p}$ | L ${ }^{1}$ | - |
| 31 |  |  | $P_{\text {c }}$ | $\mathrm{L}^{1}$ | - |
| 32 |  |  | $x_{\mathrm{P}}$ | $\mathrm{L}^{1}$ | 5 |
| 33 |  |  | $x_{\text {P, }}$ | $L^{1}$ | - |
|  | III | bounding box length maximum Feret diameter orthogonal Feret to $x_{\mathrm{Fe}, \mathrm{min}}$ |  | $\mathrm{L}^{1}$ | - |
| 35 |  |  | $x_{\mathrm{Fe}, \text { max }}$ | $\mathrm{L}^{1}$ | - |
| 36 |  |  | $x_{\text {Fe,max } 90}$ | $\mathrm{L}^{1}$ | - |
| 37 | IV | bounding box width minimum Feret diameter orthogonal Feret to $x_{\mathrm{Fe}, \max }$ | $w_{\text {bb }}$ | L ${ }^{1}$ | - |
| 38 |  |  | $x_{\mathrm{Fe} \text {,min }}$ | $\mathrm{L}^{1}$ | - |
| 39 |  |  | $x_{\text {Fe,min } 90}$ | $\mathrm{L}^{1}$ | - |
| 40 | V | min . enclosing circle diameter max. inscribed circle diameter | $d_{\text {ec }}$ | L ${ }^{1}$ | - |
| 41 |  |  | $d_{\text {ic }}$ | $L^{1}$ | - |
| 42 | VI | aspect ratio orthogonal aspect ratio | AR | - | - |
| 43 |  |  | $\mathrm{AR}_{90}$ | - | - |
| 44 | VII | circularity <br> form factor <br> Wadell's circularity bounding circles circularity | $\psi_{\mathrm{c}}$ | - | 12 |
| 45 |  |  | $\psi_{\mathrm{Kr}}$ | - | 13 |
| 46 |  |  | $\psi_{\text {Wa }}$ | - | 14 |
| 47 |  |  | $\psi$ Но | - | 15 |
| 48 | VIII | 2D solidity | $S_{x, 2 \mathrm{D}}$ | - | 17 |
| 49 |  | 2D convexity | $\mathrm{C}_{\mathrm{x}, 2 \mathrm{D}}$ | - | 19 |



Figure 9: Correlation matrices for particle characteristics determined from aligned projections, Figs. 9c through 9f only show the first quadrant (upper left) of the complete correlation matrix as shown in Fig.9a, with dashed red lines separate geometric properties and equivalent diameters from shape factors (cf. Table 2)


Figure 10: Correlation of Wadell's alternative definition for circularity (Eq. 14) and bounding circles circularity (Eq. 15)


Figure 11: Comparison of Feret diameters to 3D measures for all solids
Interestingly, elongation (but not thickness) also correlates well with Wadell's alternative circularity definition $\psi_{\mathrm{c}, \mathrm{Wa}}$ and the bounding circles
circularity $\psi_{\mathrm{bc}}$. Elongation therefore seems to be a much better indicator deviation from the cubic shape than flatness. It therefore makes sense that Krumbein takes elongation as a square in his sphericity definition, Eq. 11.

Because some of the characteristics are not correlated linearly, however, it makes sense to not stick to the Pearson correlation coefficient $r_{x y}$. Instead Fig. 9d shows the correlation matrix with the Spearman rank correlation coefficient $r_{\mathrm{s}}$. This coefficient doesn't describe a linear relationship, but rather how likely it is that a monotonic function exists between two variables. In comparison of Figs. 9c and 9d, specific surface area $S_{\mathrm{V}}(9)$ now shows a very good (Spearman) correlation with projection area $A_{\mathrm{p}}$ (1 to 4). Of course, specific surface area is directly linked to projection area, however, definitely not in a linear way, as Fig. 9b shows. Because of this advantage of finding all possible relationships instead of just the linear ones, the Spearman rank correlation coefficient is chosen for all other correlation matrices.

The mica particles stand apart from the more compact particles in several ways (Fig. 9e). Because of their flat appearance, width $(10,11)$ and length $(12,13)$ can still be approximated exceptionally well. As for the shape factors, most of the correlations are less pronounced than for the compact particles, with one exception deriving directly from the previous statement: a perfect correlation of elongation (18) and aspect ratio $(17,18)$.

The spherical soda-lime glass provides some much higher correlations for the shape factors. The 3D shape factor trifecta of elongation (18), Krumbein's (21), and Hofmann's sphericity (23) correlate highly with the 2D shape factors aspect ratio $(17,18)$, Wadell's alternative circularity $(21)$ and bounding circles circularity (22). This stresses again that the latter two circularity values correlate highly with aspect ratio, which probably diminishes their usefulness in static image analysis.

### 3.2. Stable Positioning

As described in section 2.2.1, at least the two most probable resting positions were used to produce projections. However, it is instructive to plot the distribution stable positions with a probability $p>0.1$ per solids type, as shown in Fig. 12a. Again, the soda-lime glass and mica particles clearly deviate from the compact particles (quartz, limestone, dolomite, and aluminium oxide). The compact particles on average have three to four stable positions. There are some outliers at six and even seven stable positions. One limestone particle is shown in its seven stable positions in

Fig. 12b. In contrast, the soda-lime glass spheres have no stable positions $p>0.1$ for $80 \%$ of particles. The flaky mica particles expectedly orient themselves on one of their flat sides, and so obtain on average two stable positions.


Figure 12: Stable positions of investigated solid particles for a position probability of $p>0.1$

For the simulation of static image analysis via stable positioning, the correlation matrix in Fig. 13a exhibits an expected drop in very high correlations. The 3D geometric measures of highest correlation are convex surface (6) and its equivalent diameter $x_{\mathrm{S}}(8)$; specific surface area $S_{\mathrm{V}}(9)$ scales well with projection area (1) and its equivalent diameter $x_{\mathrm{A}}$ (3); particle length $l(10,11)$ correlates highly with $x_{\mathrm{Fe}, \max }$ and $x_{\mathrm{Fe}, \max 90}$.

A few slightly more unexpected, but very high correlation values exist. Particle length $l(10,11)$ also pretty much equals the minimum enclosing circle diameter $d_{\mathrm{ec}}$ (15). Of course, the diameter in a stable position must be at least that of particle length, as this 3D dimension should always be visible in static image analysis. Only in rare cases, however, will the radius much exceed length particle length. Another high correlation is found between the minimum enclosing sphere diameter $d_{\mathrm{es}}$ (16) and the convex perimeter $P_{\mathrm{C}}(6)$ and its equivalent diameter $x_{P}$ (8), maximum Feret diameters $(10,11)$ and the minimum enclosing circle diameter $d_{\mathrm{ec}}(15)$.

As for shape factors, elongation $w / l(18)$ still correlates well with aspect ratios AR (17) and $\mathrm{AR}_{90}$ (18).

(a) Correlation coefficient matrix stable positions; correlations in red for $r_{\mathrm{S}}>0.98$, green for $r_{\mathrm{s}}>0.8$, purple for $r_{\mathrm{s}}<-0.9$

(b) Correlation matrix for random orientations; correlations in red for $r_{\mathrm{s}}>0.97$, green for $r_{\mathrm{s}}>$ 0.7 , purple for $r_{\mathrm{s}}<-0.9$

Figure 13: Correlation matrices of Spearman rank correlation coefficients for particle characteristics for compact particles; dashed red lines separate geometric properties and equivalent diameters from shape factors (cf. Table 2)

Of course, correlations between 2D and 3D particle characteristics for static image analysis, as was discussed in this and the previous section, could have been found from careful thought experiments. Wadell based his alternative sphericity definition (Eq. 14) on a projection of a particle at rest exactly because length and width should always be measurable in this situation, and most shape factors should scale will with the derived aspect ratio/elongation, as long as the particles are not deviating too much from the cubic shape.

### 3.3. Random Orientation

When comparing the correlation matrices of the stable position analysis (Fig. 13a) and that for dynamic simulation (Fig. 13b), it is first noticed that the amount of correlation is again decreasing. Note how correlations in red now have a value of $r_{\mathrm{s}} \geq 0.97$ instead of $r_{\mathrm{s}} \geq 0.98$ for the stable positions analysis.

Mostly, the properties of the 3D convex hull, $V_{\mathrm{c}}(2), x_{\mathrm{V}, \mathrm{c}}(4), S_{\mathrm{c}}(6)$, $x_{\mathrm{S}, \mathrm{c}}(8)$, and $S_{\mathrm{V}}(9)$ scale well with projection area-related characteristics $A_{\mathrm{c}}(1), x_{\mathrm{A}}(2), A_{\mathrm{c}}(3)$, and $x_{\mathrm{A}, \mathrm{c}}$ (4). Additionally, the 3D convex hull's surface area $(6,8)$ correlates well with the 2 D convex hull's perimeter $(6$, 8). However, remember that the Spearman rank correlation coefficient
is used: correlations here need not be linear. The last correlation with $r_{s} \geq 0.97$ is between specific surface area (9) and the maximum inscribed circle diameter, which is a possibly interesting starting point for further investigation.

In case of the derived shape factors, the only good correlation exists between 3D (24) and 2D solidity (23), $S_{x, 3 \mathrm{D}}$ and $S_{\mathrm{x}, 2 \mathrm{D}}$, respectively. Otherwise, shape factors do not really scale well anymore.

Especially the relationship of projection area and particle surface area is well known as Cauchy's theorem [47, 48]. Cauchy's theorem states that the surface area of a convex body $S_{p, c}$ is four times the projection area averaged over several projections $\overline{A_{\mathrm{p}, \mathrm{c}}}$.

$$
\begin{equation*}
S_{\mathrm{p}, \mathrm{c}}=4 \overline{A_{\mathrm{p}, \mathrm{c}}} \tag{20}
\end{equation*}
$$

This theorem can be tested directly on the simulated data, not so much to prove the theorem, but to test the validity of the dataset. Fig. 14 shows the relations of surface area and projection area, both for the actual particles and their convex hulls. Note that single points are plotted, not actual averaged values, so Cauchy's theorem may only hold on the average, which is why linear regression lines are included. For the compact particle convex hulls (Fig. 14b), the value of 3.92 is particularly close to the theoretical value. For both soda-lime glass and mica the values decrease. For the mica particles, the lower regression value is expected, as it is very likely for a flaky particle to produce silhouettes of comparably lower area. For the soda-lime glass spheres, the lower result may be due to the same inaccuracies of the mesh surface that lead to the maximum sphericity values of $\psi_{\mathrm{Wa}}=0.92$.

For the relation of particle surface and projection area, i.e., the nonconvex shapes, surface area overestimated for both the compact particles and the spherical soda-lime glass. This trend is no doubt because the rugged surface, but may not be unique: for high surface roughnesses, projections may underestimate actual surface area [36]. In contrast, for mica particles, surface area is still grossly underestimated because the shape effect persists.

### 3.4. Circularity vs. Sphericity

It was deemed a worthwhile exercise to see how well circularity $\psi_{c}$ and sphericity $\psi_{W a}$ correlate for the dynamic image simulation, because


Figure 14: Correlations of surface area and projection area
circularity is commonly understood as the 2D equivalent of sphericity.
A random accident led to investigation of the relationship of circularity and sphericity for the mica particles first. Fig. 15 shows the resulting correlations. The first insight is in regards to extremely small correlation values of the shape factors in the correlation matrices: at first sight, there is only a point cloud with no tendency whatsoever. At second sight, because of the nature of the two shape factors, both should be zero for infinitely stretched objects and one for spheres. Because of this unique relationship, a linear regression needs no offset, i.e., should start from zero. If a linear regression then returns a slope of one, the two shape factors are perfectly correlated. Any spread in either direction is then purely stochastic.

(a) (Wadell's) sphericity and its square root vs. circularity

(b) Sphericity and its square root vs. form factor

Figure 15: Correlation of sphericity with circularity for mica particles
From Fig. 15a it can be seen that the correlation between circularity
$\psi_{c}$ and sphericity $\psi_{W_{a}}$ is rather non-ideal, whereas the square root of sphericity $\sqrt{\psi_{W_{a}}}$ leads to a near-perfect linear regression slope of 0.99 . If this correlation is squared, we get near-perfect slope of 0.95 for sphericity $\psi_{\mathrm{Wa}}$ over the square of circularity, which is the form factor, $\psi_{\mathrm{c}}^{2}=\mathrm{FF}$.

However, the correlation of sphericity and form factor does not hold nearly as well for the compact particles. Fig. 16b shows the reslting correlation. Not only is the resulting regression slope at 1.10, but the points also do not scatter as randomly around the regression line as was the case with the mica particles.

To find if there is an underlying variable with which the data could be corrected, the data was plotted as shown in Fig. 16a. We will call Fig. 16a the parameter plot, as it shows how parameters scale within a correlation. The plots all show the same relationship, but individual points are plotted with a color map that scales according to a third parameter. To make any relationship, if existing, clear, the color map always scales between the smallest and the largest value of the chosen parameter. In the case of circularity and form factor we can see a smooth color band from left to right, which makes sense, given that the plot's $x$-axis is the form factor. To correct the point cloud to scatter more evenly around the equality line, there needs to be a parameter that changes monotonously from the upper left to the lower right of the graph, i.e., orthogonally to the equality line. Solidity, for example, is a poor candidate because it decreases in direction of the $y$-axis.

In contrast, 2D convexity $C_{x, 2 \mathrm{D}}$ fulfills the described relationship for the given data, with the smallest values found in the upper left corner, and values decreasing toward the equality line. Fig. 16b displays the same plot with a color bar for the convexity values. The parameter is thus a good candidate to correct the linear relationship of form factor and sphericity: if the form factor is divided by the 2D convexity, points in the upper left of the plot will move to the right, while points close to the equality line will stay there, as their convexity values are already close to one.

In fact, if the form factor is divided by the square of 2D convexity $C_{x, 2 D^{\prime}}^{2}$ there is, at least visually, no correlation of the data with the parameter at all anymore, as shown in Fig. 16c. However, the correlation to sphericity has worsened, with a regression slope of only 0.82 .

The procedure is thus repeated with a new parameter plot that contains the $x$ - and $y$-axes of the new correlation. The next candidate shape factor,


Figure 16: Pathway to a correlation of 2D shape factors and Wadell's sphericity; only compact particles (no soda-lime glass and mica) are shown; final result in Fig. 17
that fulfills the requirements described above is the bounding circles circularity $\psi_{\mathrm{c}, \mathrm{bc}}$, as shown in Fig. 16d. The shape factor can be used to produce an excellent correlation by "stretching" the data back to the equality line, Fig. 16e. The regression slope is now almost perfect at 1.02. Furthermore, the resulting correlation exhibits expected behavior for a correlation of circularity and sphericity: at high values approaching one, there is little error in the prediction, while the error widens as the values decrease, because there is a higher fluctuation in the projection images that can be produced for more irregular particles.

The correlation thus found is:

$$
\begin{equation*}
\psi_{\mathrm{Wa}}=\mathrm{FF} \sqrt{\psi_{\mathrm{c}, \mathrm{bc}}} / C_{\mathrm{x}, 2 \mathrm{D}}^{2} \tag{21}
\end{equation*}
$$

Several other equations were tested concerning their relevance for the given solids, i.e., for their predictive power with regards to Wadell's sphericity. The simplest correlation is the one found for mica:

$$
\begin{equation*}
\psi_{\mathrm{Wa}}=\mathrm{FF} \tag{22}
\end{equation*}
$$

Calculating the bounding circles for a given projection was a problem that was solved relatively late in this study. Because the bounding circles circularity $\psi_{c, b c}$ was therefore not available, an earlier correlation that showed the best results was identified as follows:

$$
\begin{equation*}
\psi_{\mathrm{Wa}}=\mathrm{FF} \sqrt{\mathrm{AR}_{90}} / C_{\mathrm{x}, 2 \mathrm{D}}^{2} \tag{23}
\end{equation*}
$$

Essentially, $\psi_{c, b c}$ is replaced with the orthogonal aspect ratio $\mathrm{AR}_{90}$. Given the strong correlation between the two shape factors in the static image simulations, this substitution is justified. However, because the correlation is less pronounced for dynamic image analysis, it would also be expected for Eqs. 21 and 23 to yield different results.

A combination of $\psi_{\mathrm{c}, \mathrm{bc}}$ and $\mathrm{AR}_{90}$ was also tested:

$$
\begin{equation*}
\psi_{\mathrm{Wa}}=\operatorname{FF} \psi_{\mathrm{c}, \mathrm{bc}} /\left(C_{\mathrm{x}, 2 \mathrm{D}} \mathrm{AR}_{90}\right) \tag{24}
\end{equation*}
$$

Finally, two more equations were tested so see how much the increase in number of parameters would effectively improve the correlation.

$$
\begin{equation*}
\psi_{\mathrm{Wa}}=\psi_{\mathrm{c}, \mathrm{bc}} \tag{25}
\end{equation*}
$$

$$
\begin{equation*}
\psi_{\mathrm{Wa}}=\psi_{\mathrm{c}, \mathrm{bc}} / \sqrt[3]{\mathrm{AR}_{90}} \tag{26}
\end{equation*}
$$

For all six candidate equations mentioned above, average sphericity predictions were calculated. The results are summarized in Table 3. Eq. 21 is superior compared to all others. Depending on the solid, some equation may be more accurate in their predictions. For example, Eq. 26 will give closer sphericity values for quartz and limestone, and - as expected -Eq. 22 will provide a better estimate for mica. Overall, however, Eq. 21 is the most useful generally.

Table 3: Average sphericities determined with the correlation candidates, equations 21 through 26

|  |  |  | equation |  |  |  |  |  |  |
| :--- | :---: | :---: | ---: | ---: | ---: | ---: | ---: | ---: | :---: |
| material |  | 3D | 21 | 22 | 23 | 24 | 25 | 26 |  |
| quartz | $\psi_{\mathrm{Wa}}$ | 0.71 | 0.68 | 0.63 | 0.74 | 0.62 | 0.64 | 0.70 |  |
|  | $r^{2}$ | - | 0.985 | 0.943 | 0.984 | 0.973 | 0.982 | 0.990 |  |
| limestone | $\psi_{\mathrm{Wa}}$ | 0.72 | 0.66 | 0.58 | 0.7 | 0.6 | 0.62 | 0.69 |  |
|  | $r^{2}$ | - | 0.982 | 0.928 | 0.979 | 0.970 | 0.978 | 0.990 |  |
| mica | $\psi_{\mathrm{Wa}}$ | 0.43 | 0.45 | 0.43 | 0.49 | 0.45 | 0.45 | 0.55 |  |
|  | $r^{2}$ | - | 0.824 | 0.911 | 0.854 | 0.358 | 0.219 | 0.253 |  |
| dolomite | $\psi_{\mathrm{Wa}}$ | 0.68 | 0.68 | 0.58 | 0.74 | 0.59 | 0.64 | 0.70 |  |
|  | $r^{2}$ | - | 0.986 | 0.958 | 0.985 | 0.979 | 0.982 | 0.991 |  |
| soda-lime | $\psi_{\mathrm{Wa}}$ | 0.89 | 0.93 | 0.76 | 0.94 | 0.83 | 0.91 | 0.93 |  |
|  | $r^{2}$ | - | 0.995 | 0.945 | 0.996 | 0.980 | 0.992 | 0.996 |  |
|  | $\psi_{\mathrm{Wa}}$ | 0.61 | 0.61 | 0.54 | 0.67 | 0.55 | 0.58 | 0.65 |  |
| $\mathrm{Al}_{2} \mathrm{O}_{3}$ | $r^{2}$ | - | 0.976 | 0.939 | 0.973 | 0.968 | 0.972 | 0.984 |  |
|  |  |  |  |  |  |  |  |  |  |

Fig. 17 shows the resulting correlation of Eq. 21 for all solids. As previously determined, there is significant error for soda-lime glass particles at very high sphericities due to the nature of the meshed surfaces, which results in the lowest slope of the cubic particles. For mica, the correlation is especially poor, though the average predicted sphericity is only about $5 \%$ off from the actual value.

Note that the predictive value is reasonably good because of the large number of data points. If there had been only a handful of particles, the final correlation would have been nearly impossible to find. Furthermore, the predictive power may not hold for all types of solids, especially because of the use of convexity. If surface roughness significantly increases,


Figure 17: Best correlation of 2D shape factors and Wadell's sphericity
surface area effects could be underestimated by 2D convexity. Because the resolution of STL mesh, voxel image, and projection silhouette are directly linked and should be identical, the correlation is expected to give sphericity values at the same resolution for the surface area of the particle.

## 4. Conclusions

A collection of particle surface meshes, resulting from X-ray tomographic measurements, has been used to simulate both static and dynamic image analysis. The results have been evaluated to find the highest correlations between 2D and 3D geometric measures and shape factors. The dataset and methods described prove to be physically accurate, although highly spherical soda-lime glass particles reach a final sphericity lower than one because of the nature of the description of particle surfaces as triangular meshes.

A correlation between Wadell's sphericity in 3D and the form factor in 2D has been found that is expected to predict sphericity values well for a wide range of particles, provided that enough particles are measured. Confirmation experiments with a broader set of particles are planned in the future.

The dataset, as provided in the supplementary data, offers the possibility to discover numerous correlations and insights regarding geometric measures and shape factors, as well as their relationships across two and
three dimensions. We encourage researchers to use the dataset for their research questions and to shed light into questions that had long been obscured by computational complexity.

## Supplementary Data

Supplementary files are available in the Open Access Repository and Archive for Research Data of Saxon Universities (OPARA):
https://doi.org/10.25532/OPARA-479
Supplementary files enable users to reproduce imaging datasets as used in this study and demonstrate the methods for acquisition of all particle characteristics for an example particle. Particle STL files and the resulting dataset tables are included. Note that you need a working Python setup and that all code is made available as Jupyter notebooks.

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