

Thermal Bianisotropy

Gal Shmuel*

Department of Mechanical Engineering, Technion-Israel Institute of Technology, Israel

John R. Willis†

Centre for Mathematical Sciences, Cambridge University

(Dated: May 26, 2025)

Breaking spatial symmetries can induce interactions between disparate physical fields, which manifest in the macroscopic properties of materials as cross-coupling terms. Prominent examples include Willis terms in phononics and bianisotropic terms in photonics. However, the development of analogous cross-couplings in heat conduction remains limited and incomplete; here, we close this knowledge gap. To this end, we introduce an exact, universal method for describing the macroscopic dynamics of various physical processes. For heat conduction and thermodynamics, the method shows that thermal bianisotropy emerges through the intentional design of spatial asymmetry. The resulting macroscopic description eliminates the infinite heat speed paradox inherent in Fourier’s law of conduction. We support the exact theory with examples based on heuristic homogenization of a canonical scattering problem. These results may benefit the design of thermal metamaterials with asymmetric response.

Metamaterials are composite materials architected to exhibit extraordinary effective properties, either in value or in nature [1]. A route to calculate and design such properties is via homogenization theories, aimed at extracting the macroscopic features of materials with microstructure. Willis [2–5] developed an exact homogenization theory for elastodynamics, revealing several metamaterial features, perhaps the most notable being the unusual macroscopic coupling between momentum and strain. The advent of metamaterials revived interest in this theory [6–11], mainly since these cross-couplings provide additional mechanisms for wave manipulation [12–18].

The applicability of the homogenization framework developed by Willis goes beyond elastodynamics. In electromagnetics, it predicts bianisotropic response, even when constituent materials have no magnetoelectric couplings [19], a conclusion that was obtained by the electromagnetics community using other means [20–22]). A generalization of the theory by Shmuel and Pernas-Salomón [23] to piezoelectric media led to the discovery of metamaterial coupling between momentum and electric field [24–28], earning the name ‘electromomentum coupling’.

Here, we adapt these homogenization schemes to heat conduction and thermodynamics. This adaptation is not trivial, owing to the difference in the equations that govern heat conduction relative to other transport phenomena. Specifically, classical heat conduction is governed a parabolic-type partial differential equation, whereas the type of equations that govern elastodynamics, electromagnetics and piezoelectricity is hyperbolic. Nevertheless, we find that the unusual cross-couplings appear also in the macroscopic response of thermal media, now between heat flux and temperature; and between entropy and temperature gradient.

A closely related coupling was previously introduced by Torrent et al. [29] and Xu et al. [30], who coined the name

thermal Willis coupling. Our conclusions are different and substantially broader, as we detail next. To do so, we recall that Ref. [30] begins with the statement “*Willis coupling generically stems from bianisotropy or chirality in wave systems. Nevertheless, those schemes are naturally unavailable in diffusion systems described by a single constitutive relation governed by the Fourier law*”. With this belief, to realize thermal Willis coupling Refs. [29, 30] modulated the constitutive properties of the medium both in time and space, which is an extremely challenging design task. However, we prove that spatial asymmetry is sufficient and employ Willis’ homogenization schemes originally developed for elastodynamics. In doing so, we interpret the relation between entropy and temperature as a constitutive equation; it is this constitutive equation that was missing in Refs. [29, 30]. As a result, we also provide an expression for the mean entropy that is absent in Refs. [29, 30]. There is also a difference in the macroscopic heat equation in Ref. [29] and in Ref. [30]. This is an outcome of ambiguity in defining the nonlocal effective relations within their Bloch-Floquet framework. Simply put, since spatial and temporal derivatives become products with wavenumber and frequency, respectively, they are mixed up in their transformed domain [8, 24]. As a result, there are different possibilities to invert the effective relations back into the spatiotemporal domain. Following Fietz and Shvets [31] and Willis [19], here we eliminate this ambiguity by introducing a residual field which acts as additional source [10] and renders the mean field variables independent. Finally, our homogenization is broader in that it applies to three-dimensional media which may be periodic or random, finite or infinite, while the homogenization schemes in Refs. [29, 30] are limited to infinite one-dimensional media modulated in one direction.

Before we proceed to the analysis, we highlight a major implication of our homogenization scheme: in effect, the nature of the heat equation for certain composites is no longer parabolic and terms that are nonlocal in space and time now dictate thermal relaxation. In the spatially-local approximation, examples show that these terms eliminate the unrealistic prediction of infinite speed of thermal signals in classical heat

* meshmuel@technion.ac.il

† jrw1005@cam.ac.uk

conduction theory. Formulating heat conduction theories with finite speed [32] (so-called second sound [33]) is the focus numerous works which date back to Cattaneo's pioneering paper [34, 35]. Notably, we arrive at a reminiscent result without postulating from the outset a dependency on the history of the temperature; rather, it is an inevitable outcome of our homogenization scheme.

We begin by recalling the equations that govern heat conduction. Using Gibbs relations, we write the energy balance equation as

$$\theta_R \dot{\eta} + \nabla \cdot \mathbf{q} = r, \quad (1)$$

where η and r are the entropy and heat source per unit volume, respectively, \mathbf{q} is the heat flux per unit area and θ_R is the reference temperature. Fourier's law, $\mathbf{q} = -\kappa \nabla \theta$, states that \mathbf{q} is coupled with the gradient temperature change $\nabla \theta$ by the thermal conductivity tensor κ . In addition, the temperature and entropy are coupled by the specific heat c such that $\theta_R \dot{\eta} = c \dot{\theta}$. Recognizing the latter as a constitutive equation is key in adapting Willis' homogenization method [19] from electromagnetics and elastodynamics to heat conduction. We provide the complete details of the scheme in a companion paper [36], and summarize here its key principles and main results. The method uses ensemble averaging in order to define the effective fields, such that they identically satisfy the macroscopic equations. The method also employs a residual kinematic field which delivers a unique definition of effective relations. After some work using Green's functions and adjoint equations, we are able to express the effective heat flux and entropy in terms of the effective temperature and its gradient. These effective relations are nonlocal in space and time and take bianisotropic form

$$\begin{pmatrix} -\langle \mathbf{q} \rangle \\ \langle \theta_R \dot{\eta} \rangle \end{pmatrix} = \begin{pmatrix} \tilde{\kappa} & \tilde{\chi} \\ \tilde{\xi} & \tilde{c} \end{pmatrix} \begin{pmatrix} \langle \nabla \theta - \zeta \rangle \\ \langle \dot{\theta} \rangle \end{pmatrix}, \quad (2)$$

where $\langle \cdot \rangle$ denotes ensemble averaging, ζ is the residual temperature gradient and $\tilde{\chi}$ and $\tilde{\xi}$ are bianisotropic vector terms that are absent from the microscopic constitutive relations. The structure of the kernel in Eq. (2) suggests the following alternative form

$$\begin{pmatrix} -\langle \mathbf{q} \rangle \\ \langle \theta_R \dot{\eta} \rangle \end{pmatrix} = \begin{pmatrix} \langle \kappa \rangle & \mathbf{0} \\ \mathbf{0} & \langle c \rangle \end{pmatrix} \begin{pmatrix} \langle \nabla \theta - \zeta \rangle \\ \langle \dot{\theta} \rangle \end{pmatrix} + \begin{pmatrix} \hat{\kappa} & \hat{\xi}^\dagger \\ \hat{\xi} & \hat{c} \end{pmatrix} \begin{pmatrix} \langle \nabla \theta - \zeta \rangle \\ \langle \dot{\theta} \rangle \end{pmatrix}, \quad (3)$$

where $\hat{\xi}$ and $\hat{\xi}^\dagger$ are an adjoint pair if the microscopic relations are self-adjoint. To demonstrate the implication of Eq. (2) in the simplest way, we substitute its local limit in the one-dimensional case into the macroscopic heat equation, and then set $\zeta = 0$. We obtain

$$\tilde{\kappa} \partial_x^2 \langle \theta \rangle = \hat{c} \partial_t^2 \langle \theta \rangle + \langle c \rangle \partial_t \langle \theta \rangle + (\hat{\xi} - \hat{\xi}^\dagger) \partial_{xt} \langle \theta \rangle. \quad (4)$$

In some of the examples that we provide $\hat{c} > 0$, thus implying finite the speed of macroscopic thermal disturbances. In

the local limit of a reciprocal solid, the contributions of the bianisotropic terms cancel each other out in the wave equation since $\hat{\xi} = \hat{\xi}^\dagger$, and the equation takes the form of the Telegrapher's equation. Nevertheless, the cross-couplings still modify the thermal impedance $Z = \theta/q$ and thus break the symmetry of heat conduction. Specifically, the characteristic thermal impedance in the $\pm x$ direction in the presence of a heat source with time-dependency $e^{i\omega t}$ is

$$Z^\pm = (\mp \tilde{k} \tilde{\kappa} - i\omega \tilde{\xi})^{-1}, \quad (5)$$

where $\tilde{k} = \sqrt{i\omega \tilde{c} / \tilde{\kappa}}$ and evidently $\tilde{\xi}$ introduces a direction-dependent phase shift.

Motivated by the insights gained by the exact homogenization method, we proceed to quantify this symmetry breaking in a 1D scattering problem and develop a heuristic homogenization scheme in the long-wavelength limit. To this end, we consider a setting in which the heterogeneous medium acts as a thermal scatterer in-between two homogeneous semi-infinite slabs with specific heat c_0 and thermal conductivity κ_0 (Fig. 1). The slabs support leftward- and rightward evanescent solutions of the form $\theta(x, t) = \Theta(x) e^{i\omega t}$, where

$$\Theta(x) = \begin{cases} a_L^- e^{k_0(x+l/2)} + a_L^+ e^{-k_0(x-l/2)}, & x \leq -l/2, \\ a_R^- e^{k_0(x+l/2)} + a_R^+ e^{-k_0(x-l/2)}, & x \geq l/2, \end{cases}, \quad (6)$$

$k_0 = \sqrt{i\omega c_0 / \kappa_0}$ is the wavenumber and some of the coefficients $\{a_{L,R}^\pm\}$ may vanish, depending on the excitation. Specifically, if the only heat source is within the finite medium, then a_L^+ and a_R^- must vanish for the temperature to be finite at $\pm\infty$. However, if heat sources exist in the semi-infinite slabs, a_L^+ and a_R^- admit nonzero values owing to incoming heat from the sources to the scatterer. In any case, the continuity of the temperature and heat flux at the interfaces between the slabs and the scatterer couples $\mathbf{a}_L = (a_L^+, a_L^-)^T$ and $\mathbf{a}_R = (a_R^+, a_R^-)^T$ by the scattering matrix \mathbf{K} . We detail the derivation of its components in the companion paper [36], noting here that \mathbf{K} has the form

$$\mathbf{K} = \begin{pmatrix} t_L - r_R t_R^{-1} & r_L t_R^{-1} \\ -r_L t_R^{-1} & t_R^{-1} \end{pmatrix}, \quad (7)$$

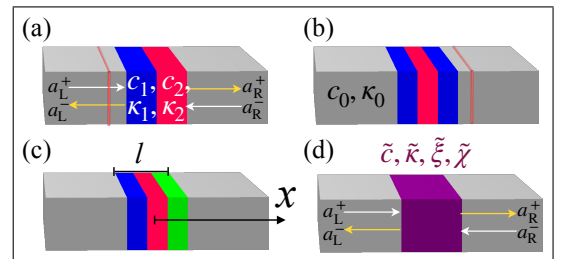


Figure 1. Heterogeneous scatterers in-between two homogeneous semi-infinite slabs. Red lines in panels (a) and (b) denote planar heat sources at the left- and right slabs (d). Homogeneous scattering with effective bianisotropic terms that generates that same scattering response as the heterogeneous scatterer.

such that r_L and t_L (r_R and t_R) are the reflection and transmission coefficients owing to a source in the left (right) slab, respectively. We find that $t_L = t_R$ as reciprocity requires, however $r_L \neq r_R$; it is this asymmetry that is captured by the effective bianisotropic terms, as we show in the numerical examples to follow.

Our heuristic homogenization aims at replacing the heterogeneous scatterer by a fictitious homogeneous medium that generates the same measurable fields, that is the same scattering matrix K . At the outset, we note that since K has three independent components (r_L, r_R and $t_L \equiv t_R$), any attempt to reproduce it using standard constitutive relations with only two material parameters (c and κ) will inevitably lead to inconsistencies. To proceed with the homogenization process, we derive the transfer matrix that relates the temperature and heat flux at both ends of a fictitious bianisotropic medium, which is governed by the local form of Eq. (2). We obtain the following expressions for the effective properties by requiring that the transfer matrix of the bianisotropic medium matches that of the heterogeneous scatterer:

$$\tilde{\kappa} = M_{12}^{-1}, \tilde{c} = -\frac{\det M}{i\omega M_{12}}, \tilde{\xi} = \frac{M_{22}}{i\omega M_{12}}, \tilde{\chi} = -\frac{M_{11}}{M_{12}}, \quad (8)$$

where the transfer matrix of the heterogeneous scatterer is $T = e^{Ml}$. Since $\det T = 1$, it follows that $\text{tr} M = 0$ and hence $\tilde{\chi} = i\omega\tilde{\xi}$.

We proceed to evaluate numerical examples, and begin with a bilayer made of SiO₂ and diamond, each layer having thickness 25 nm [Fig. 1(a)]. Fig. 2 shows the effective properties in the frequency domain calculated with Eq. (8). We observe that $\tilde{\kappa}$ and \tilde{c} are complex-valued functions that recover the harmonic- and arithmetic means, respectively, in the low-frequency limit. Both the real- and imaginary parts of these effective functions increase with frequency. The parameter \tilde{c} is obtained by subtracting from \tilde{c} the arithmetic mean of c . The real part of \tilde{c} is an even function. It is of order ω^2 as $\omega \rightarrow 0$, reducing to a linear dependence at larger frequency. It contributes a modest amount of dispersion. The imaginary part of \tilde{c} is very closely proportional to $i\omega$, corresponding to the second derivative of $\langle\theta\rangle$ in Eq. (4). Importantly, $\tilde{\chi}$ is purely imaginary and equal to $i\omega\tilde{\xi}$. For the entire range of frequency plotted, $\tilde{\xi}$ is constant and thus $\tilde{\chi}$ is linear in ω .

The second example considers symmetric trilayer made of

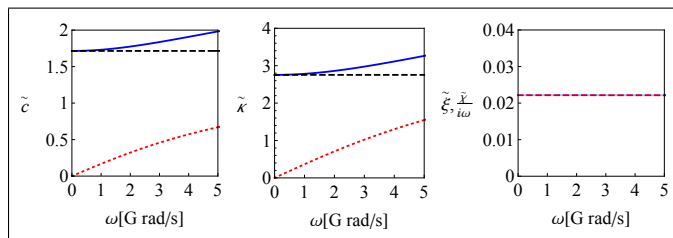


Figure 2. Effective (a) specific heat and (b) thermal conductivity properties of a bilayer made of SiO₂ and diamond. Solid blue and dashed red curves denote real and imaginary parts, respectively. Dashed black lines denote the harmonic and arithmetic means. (c) Effective bianisotropic terms $\tilde{\xi}$ (dashed red) and $\tilde{\chi}/i\omega$ (solid blue).

a diamond layer in-between two SiO₂ layers, where the thickness of each layer is 16.67 nm [Fig. 1(b)]. Fig. 3 mirrors Fig. 2, now for the trilayer. Here again, $\tilde{\kappa}$ and \tilde{c} are complex functions which recover the harmonic and arithmetic means, respectively, in the low-frequency limit. However, the imaginary part of \tilde{c} is negative at small ω . It displays a minimum around 3 G rad/s and thereafter increases. We have confirmed that this increase is linear in ω over a substantial range beyond the minimum, which is sufficient to deliver to a good approximation of the second derivative term in Eq. (4), with the lower frequencies contributing to the dissipation. Notably, $\tilde{\chi}$ and $\tilde{\xi}$ are identically zero, as expected due to the symmetry of the scatterer. The effective properties in Fig. 4 are obtained by replacing the material of the left layer with copper [Fig. 1(c)], a symmetry breaking that yields nonzero $\tilde{\chi}$ and $\tilde{\xi}$, like in Fig. 2. Cross-coupling terms of the asymmetric trilayer are greater than those of the bilayer, demonstrating that it is possible to optimize thermal bianisotropy by microstructure design.

The magnitude of $\tilde{\xi}$ is directly related to the asymmetry in the reflection properties. To illustrate this relation, we plot in Fig. 5 the magnitude ratio and phase difference between the reflection coefficients (r_L and r_R) as functions of frequency, for the different scatterers shown Figs. 2-4. We observe that a larger magnitude of $\tilde{\xi}$ corresponds to a greater relative phase between r_L and r_R .

To conclude, we have derived a formal expression for the effective response of heat-conducting composites, showing that this response inevitably must display thermal bianisotropy, and have supported this observation with simple examples based on heuristic homogenization of subwavelength elements in canonical gedanken scattering experiments. Our demonstration of thermal bianisotropy opens avenues for novel designs of thermal metamaterials, akin to how Willis materials have expanded the design space for elastic and electroelastic metamaterials. It also sets the stage for our future work on thermomomentum Willis-type couplings in thermoelastic materials, mirroring the the electromomentum effect in piezoelectric materials.

G.S. thanks Prof. Oleg Gendelman for fruitful discussions, and funding by the European Union (ERC, EXCEPTIONAL, Project No. 101045494). Funded by the European Union. Views and opinions expressed are however those of the author(s) only and do not necessarily reflect those of the European Union or the European Research Council Executive

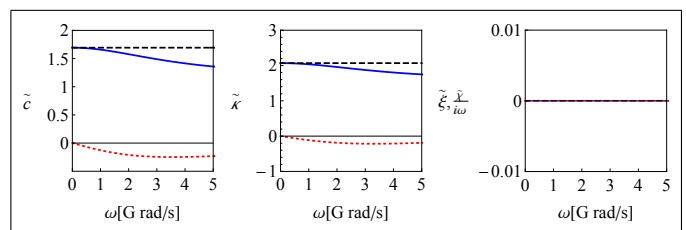


Figure 3. Effective properties of a symmetric trilayer made of a diamond layer in-between two SiO₂ layers. Legend same as in Fig. 2. The bianisotropic terms are null, owing to the symmetry of the scatterer.

Agency. Neither the European Union nor the granting authority can be held responsible for them.

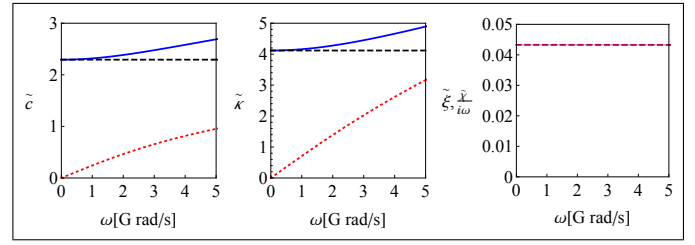


Figure 4. Effective properties of a trilayer made of SiO₂, diamond and copper layers. Legend same as in Fig. 2.

- [1] Muamer Kadic, Graeme W. Milton, Martin van Hecke, and Martin Wegener. 3d metamaterials. *Nature Reviews Physics*, 1(3):198–210, 2019.
- [2] J R Willis. Variational principles for dynamic problems for inhomogeneous elastic media. *Wave Motion*, 3(1):1–11, 1981.
- [3] J R Willis. The nonlocal influence of density variations in a composite. *Int. J. Solids Struct.*, 21(7):805–817, 1985.
- [4] J R Willis. Continuum Micromechanics. chapter Dynamics of Composites, pages 265–290. Springer-Verlag New York, Inc., New York, NY, USA, 1997.
- [5] John R. Willis. The construction of effective relations for waves in a composite. *Comptes Rendus Mécanique*, 340(4):181 – 192, 2012. Recent Advances in Micromechanics of Materials.
- [6] Michael B Muhlestein, Caleb F Sieck, Andrea Alù, and Michael R Haberman. Reciprocity, passivity and causality in Willis materials. *Proc. R. Soc. London A Math. Phys. Eng. Sci.*, 472(2194), 2016.
- [7] A. L. Shuvalov, A. A. Kutsenko, A. N. Norris, and O Poncelet. Effective Willis constitutive equations for periodically stratified anisotropic elastic media. *Proc. R. Soc. London A Math. Phys. Eng. Sci.*, 467(2130):1749–1769, 2011.
- [8] Hussein Nassar, Qi-Chang He, and Nicolas Auffray. Willis elastodynamic homogenization theory revisited for periodic media. *Journal of the Mechanics and Physics of Solids*, 2015.
- [9] H Nassar, Q.-C. He, and N Auffray. On asymptotic elastodynamic homogenization approaches for periodic media. *J. Mech. Phys. Solids*, 88:274–290, 2016.
- [10] Caleb F Sieck, Andrea Alù, and Michael R Haberman. Origins of Willis coupling and acoustic bianisotropy in acoustic metamaterials through source-driven homogenization. *Phys. Rev. B*, 96(10):104303, 2017.
- [11] Shixu Meng and Bojan B Guzina. On the dynamic homogenization of periodic media: Willis’ approach versus two-scale paradigm. *Proc. R. Soc. London A Math. Phys. Eng. Sci.*, 474(2213), 2018.
- [12] Steven R. Craig, Bohan Wang, Xiaoshi Su, Debasish Banerjee, Phoebe J. Welch, Mighten C. Yip, Yuhang Hu, and Chengzhi Shi. Extreme material parameters accessible by active acoustic metamaterials with Willis coupling. *The Journal of the Acoustical Society of America*, 151(3):1722–1729, 03 2022.
- [13] Joshua Lau, Suet To Tang, Min Yang, and Zhiyu Yang. Coupled Decorated Membrane Resonators with Large Willis Coupling. *Phys. Rev. Applied*, 12:014032, Jul 2019.
- [14] Aurélien Merkel, Vicent Romero-García, Jean-Philippe Groby, Jensen Li, and Johan Christensen. Unidirectional zero sonic reflection in passive \mathcal{PT} -symmetric Willis media. *Phys. Rev. B*, 98:201102, Nov 2018.
- [15] Michael B. Muhlestein, Caleb F. Sieck, Preston S. Wilson, and Michael R. Haberman. Experimental evidence of Willis coupling in a one-dimensional effective material element. *Nature Communications*, 8:15625 EP –, 06 2017.
- [16] Li Quan, Younes Ra’di, Dimitrios L Sounas, and Andrea Alù. Maximum Willis Coupling in Acoustic Scatterers. *Phys. Rev. Lett.*, 120(25):254301, 2018.
- [17] Junfei Li, Ailing Song, and Steven A. Cummer. Bianisotropic acoustic metasurface for surface-wave-enhanced wavefront transformation. *Phys. Rev. Applied*, 14:044012, Oct 2020.
- [18] Xiaoshi Su and Andrew N. Norris. Retrieval method for the bianisotropic polarizability tensor of Willis acoustic scatterers. *Phys. Rev. B*, 98:174305, Nov 2018.
- [19] J R Willis. Effective constitutive relations for waves in composites and metamaterials. *Proc. R. Soc. London A Math. Phys. Eng. Sci.*, 467(2131):1865–1879, 2011.
- [20] Ricardo Marqués, Francisco Medina, and Rachid Rafii-El-Idrissi. Role of bianisotropy in negative permeability and left-handed metamaterials. *Phys. Rev. B*, 65:144440, Apr 2002.
- [21] Andrea Alù. First-principles homogenization theory for periodic metamaterials. *Phys. Rev. B*, 84:075153, Aug 2011.
- [22] Andrea Alù. Restoring the physical meaning of metamaterial constitutive parameters. *Phys. Rev. B*, 83:081102, Feb 2011.
- [23] René Pernas-Salomón and Gal Shmuel. Symmetry breaking creates electro-momentum coupling in piezoelectric metamaterials. *Journal of the Mechanics and Physics of Solids*, 134:103770, 2020.
- [24] René Pernas-Salomón and Gal Shmuel. Fundamental principles for generalized Willis metamaterials. *Phys. Rev. Applied*, 14:064005, Dec 2020.
- [25] René Pernas-Salomón, Michael R. Haberman, Andrew N. Nor-

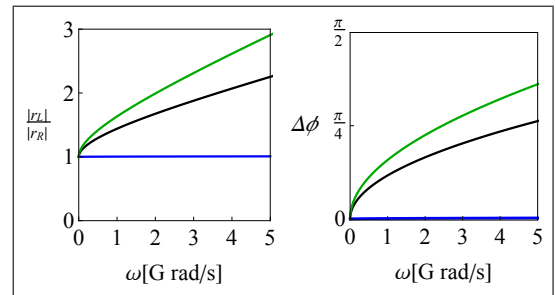


Figure 5. Magnitude ratio (left panel) and phase difference (right panel curves) between the reflection coefficients (r_L and r_R) as functions of frequency, for the different scatterers shown in Figs. 2 (black), 3 (blue) and 4 (green).

- ris, and Gal Shmuel. The electromomentum effect in piezoelectric Willis scatterers. *Wave Motion*, page 102797, 2021.
- [26] Alan Muhafra, Majd Kosta, Daniel Torrent, René Pernas-Salomón, and Gal Shmuel. Homogenization of piezoelectric planar Willis materials undergoing antiplane shear. *Wave Motion*, 108:102833, 2022.
- [27] Majd Kosta, Alan Muhafra, Rene Pernas-Salómon, Gal Shmuel, and Oded Amir. Maximizing the electromomentum coupling in piezoelectric laminates. *International Journal of Solids and Structures*, 254-255:111909, 2022.
- [28] Kevin Muhafra, Michael R. Haberman, and Gal Shmuel. Discrete one-dimensional models for the electromomentum coupling. *Phys. Rev. Appl.*, 20:014042, Jul 2023.
- [29] Daniel Torrent, Olivier Poncelet, and Jean-Christophe Batsale. Nonreciprocal thermal material by spatiotemporal modulation. *Phys. Rev. Lett.*, 120:125501, Mar 2018.
- [30] Liujun Xu, Guoqiang Xu, Jiabin Li, Ying Li, Jiping Huang, and Cheng-Wei Qiu. Thermal Willis coupling in spatiotemporal diffusive metamaterials. *Phys. Rev. Lett.*, 129:155901, Oct 2022.
- [31] Chris Fietz and Gennady Shvets. Current-driven metamaterial homogenization. *Physica B: Condensed Matter*, 405(14):2930 – 2934, 2010. Proceedings of the Eighth International Conference on Electrical Transport and Optical Properties of Inhomogeneous Media.
- [32] Morton E. Gurtin and A. C. Pipkin. A general theory of heat conduction with finite wave speeds. *Archive for Rational Mechanics and Analysis*, 31(2):113–126, 1968.
- [33] D. S. Chandrasekharaiah. Thermoelasticity with Second Sound: A Review. *Applied Mechanics Reviews*, 39(3):355–376, 03 1986.
- [34] Carlo Cattaneo. Sulla conduzione del calore. *Atti Sem. Mat. Fis. Univ. Modena*, 3:83–101, 1948.
- [35] Gianfranco Capriz, Krzysztof Wilmanski, and Paolo Maria Mariano. Exact and approximate maxwell-cattaneo-type descriptions of heat conduction: A comparative analysis. *International Journal of Heat and Mass Transfer*, 175:121362, 2021.
- [36] Gal Shmuel and John R. Willis. An exact homogenization method in heat conduction. *To be submitted to Phys. Rev. Appl.*